

EXPERIMENTAL OBSERVATION OF INTRINSIC LOCALIZED MODES IN GERMANIUM

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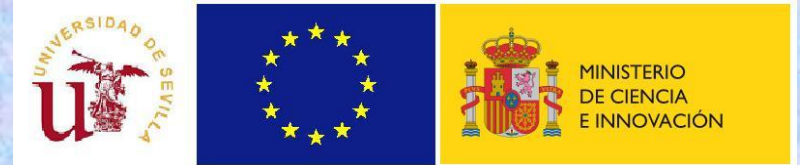
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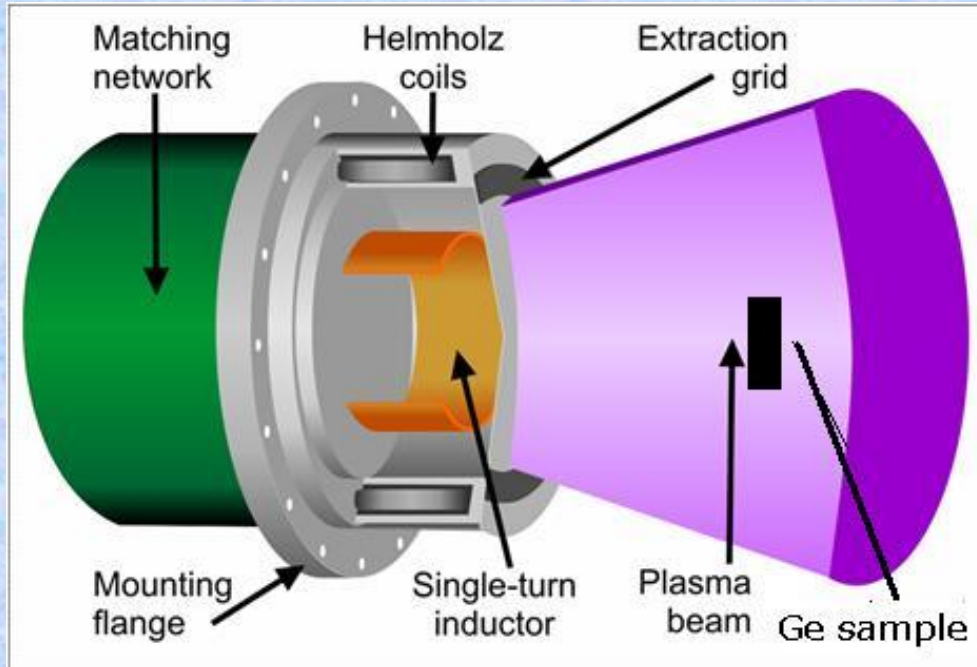


V International Symposium on Strong Nonlinear Vibronic and Electronic Interactions in Solids

Tartu, Estonia, May 1-3, 2015

Our basic experiment: Plasma annealing of defects

Copra ICP Ar plasma source



Sb-doped Ge

Defect detection by
DLTS



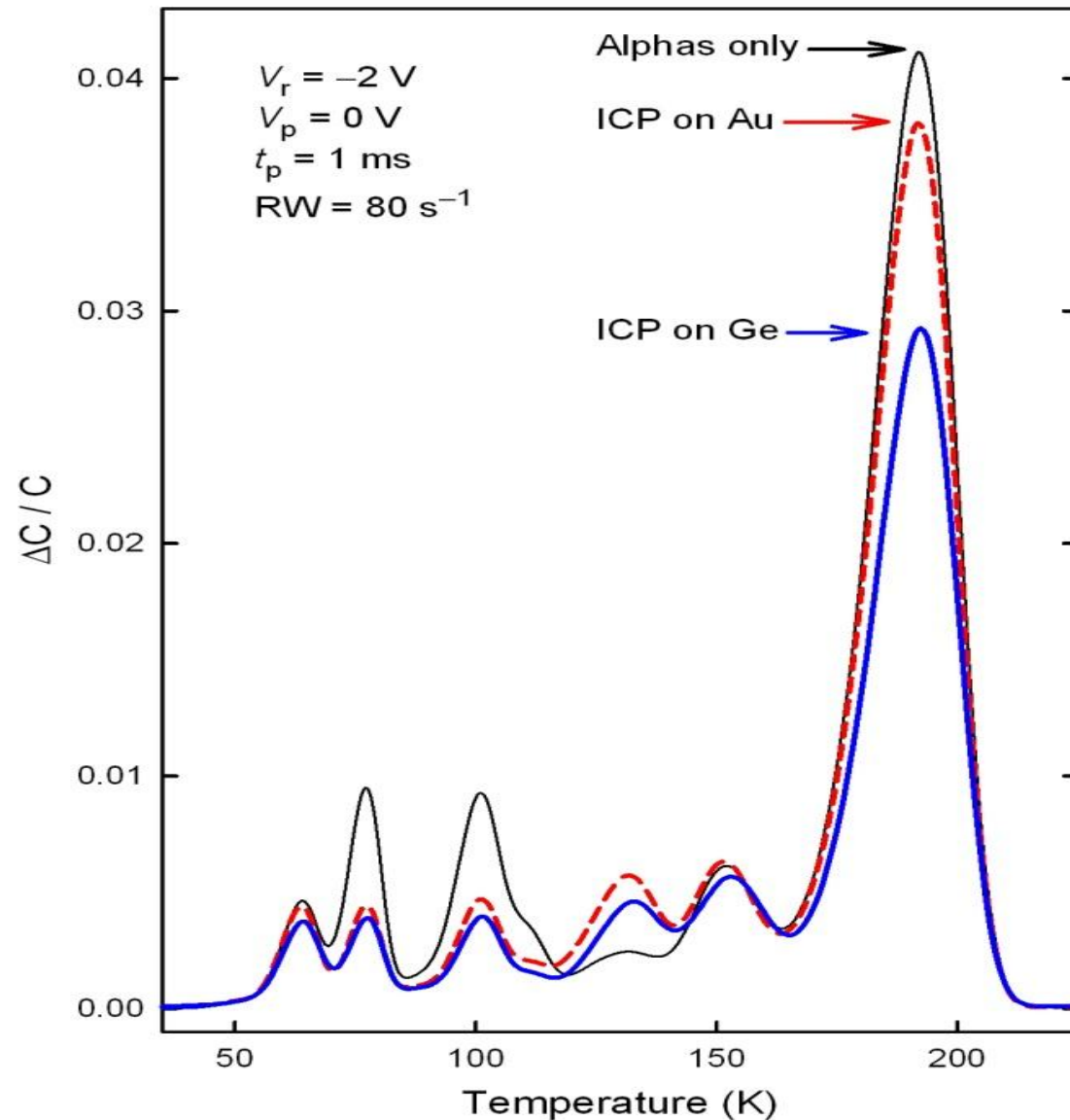
Ge sample with an Au diode

$$N_{\text{Ge}} = 4.42 \times 10^{22} \text{ cm}^{-3}$$

$$N_{\text{Sb}} = 1.03 \times 10^{15} \text{ cm}^{-3}$$

$$N_{\text{T}} = 1.07 \times 10^{14} \text{ cm}^{-3}$$

Our basic experiment: 30m, 4 eV ICP plasma annealing



Defects by alpha particles

Deep into Ge 2.6 microns

Reduction of defects in 30%

Similar rate of thermal annealing
at 150 C needs 10^{16} more energy

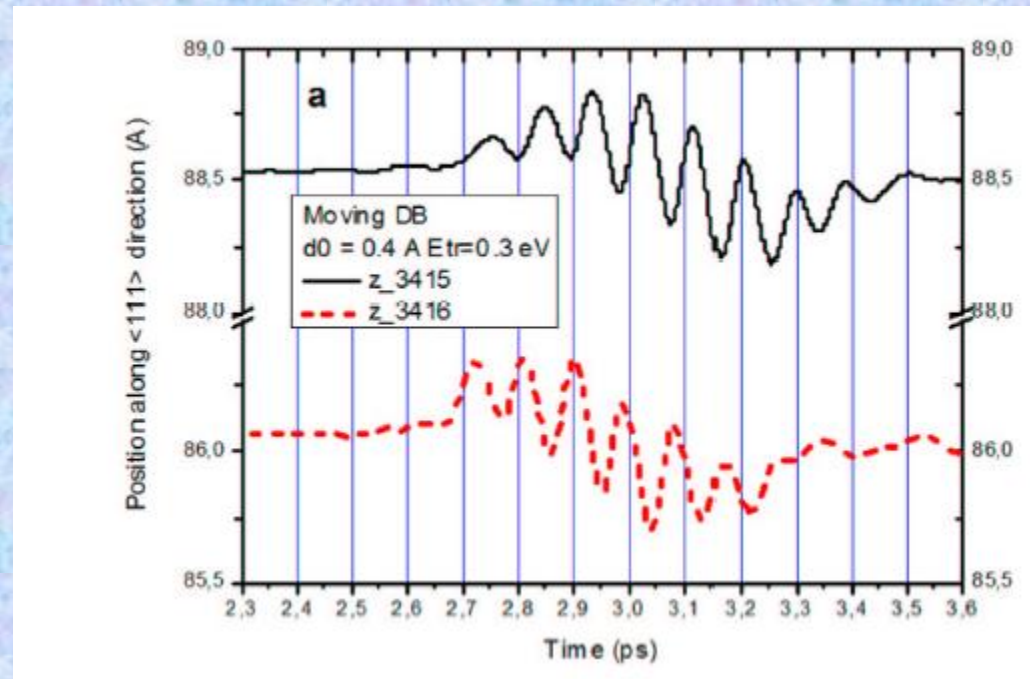
Defect detection by DLTS

Reduction through Au diode
(25nm) 7%

Very low flux

ILMs: Intrinsic (nonlinear) localized modes

- Moving ILMs
 - A perturbation of a medium that transport energy without mass transport
 - In a localized way
 - With little dispersion
 - It has a nonlinear vibrational component
 - It may have an electrical or other type component

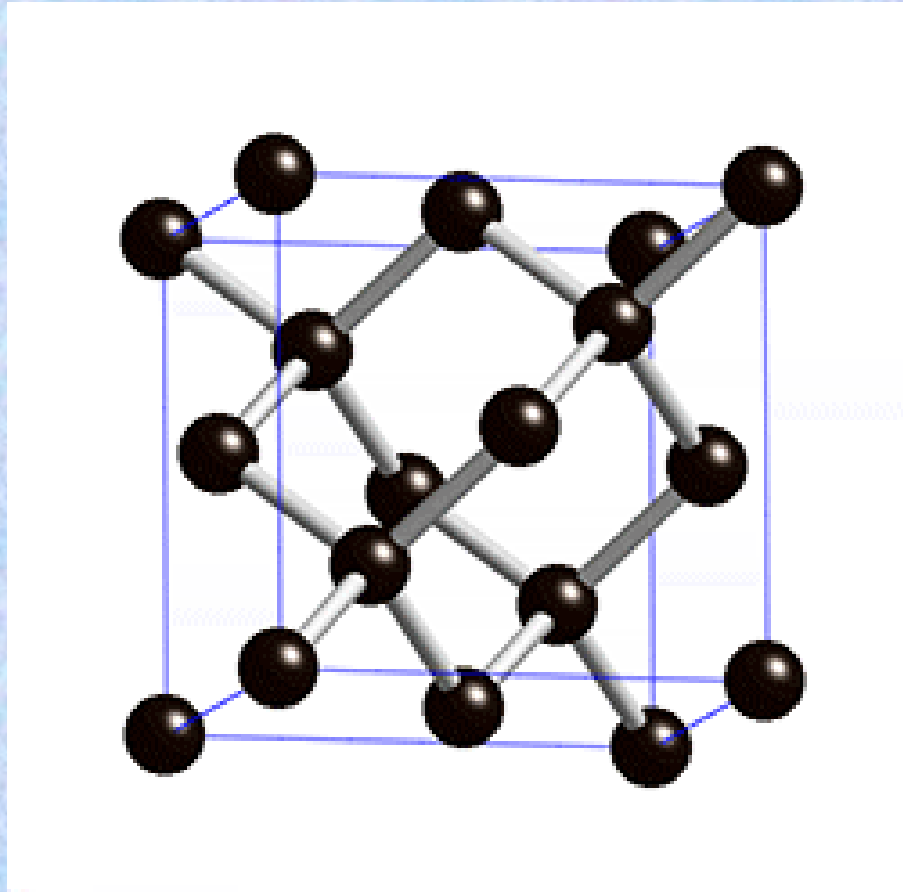


Terentyev, Dubinko, Dubinko, Dmitriev, Zhurkin, 2015

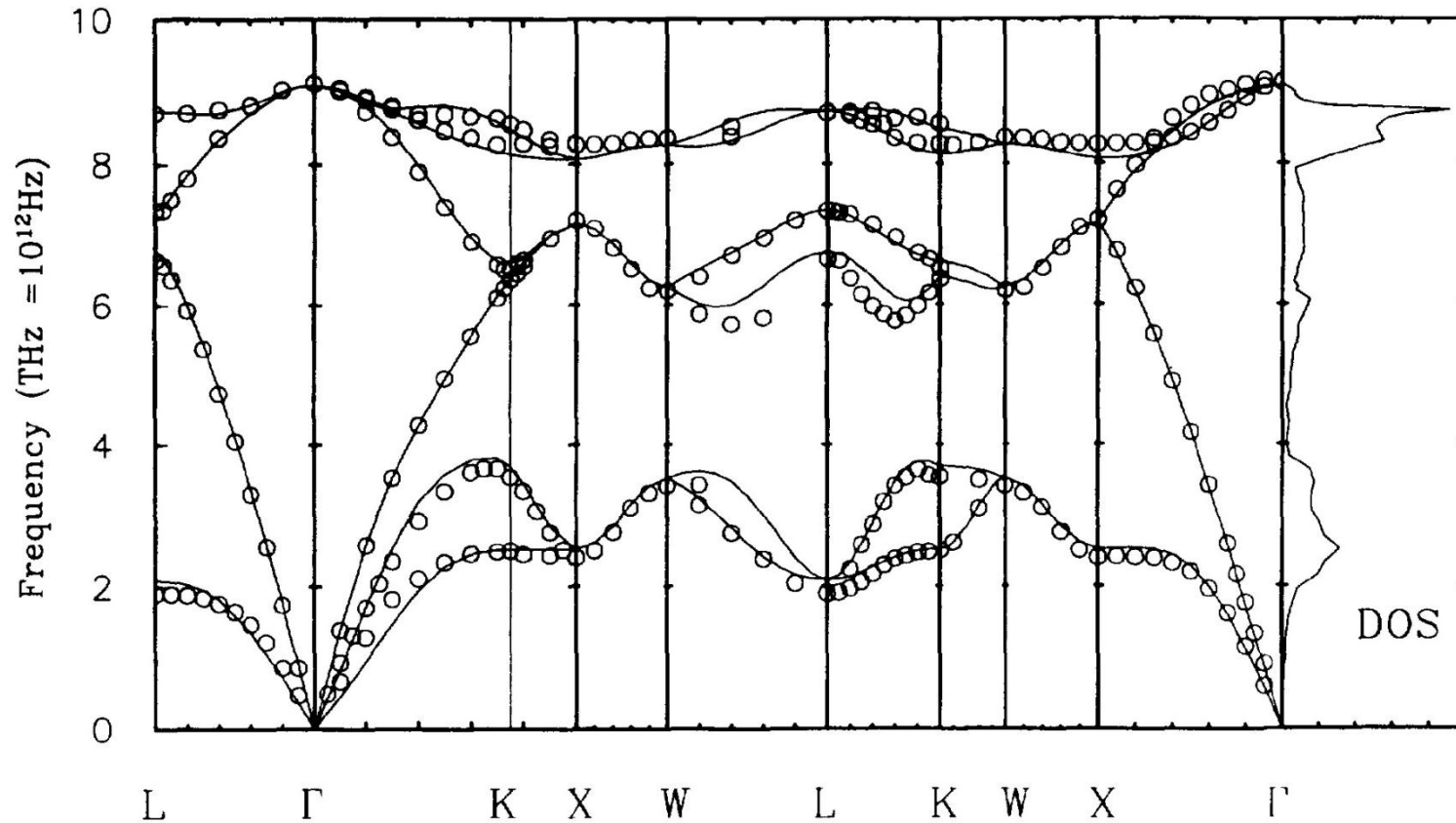
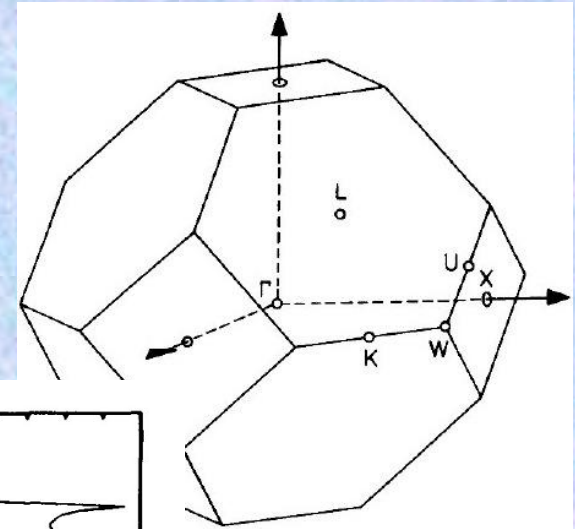
- Standing ILM: localized energy that keeps localized, mainly vibrational

ILMs in Ge?

Stationary ones obtained by Voulgarakis et al (2004)
and Hiznyakov et al (2014)



Where can be ILMs in Ge?



ILMs: how we can detect them?

- How many they are?

They are very few $\sim 10^{-9}$ number of phonons

- What energy they have? ~ 1 eV

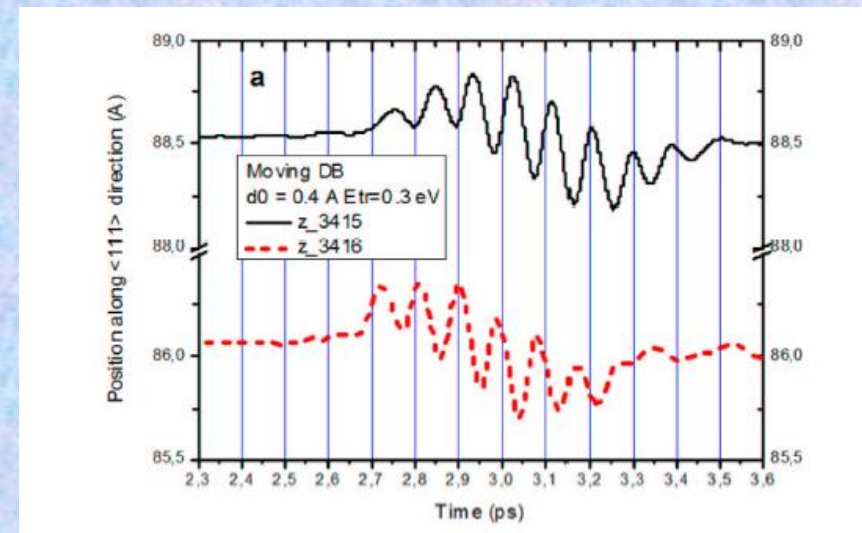
- How to detect them?

What property they change?

Something $\frac{dN}{dt} \propto \exp\left(-\frac{E_A}{k_B T}\right)$

- How far they travel?

100-1000 lattice units?



ILM detection: amplification effect of Arrhenius law

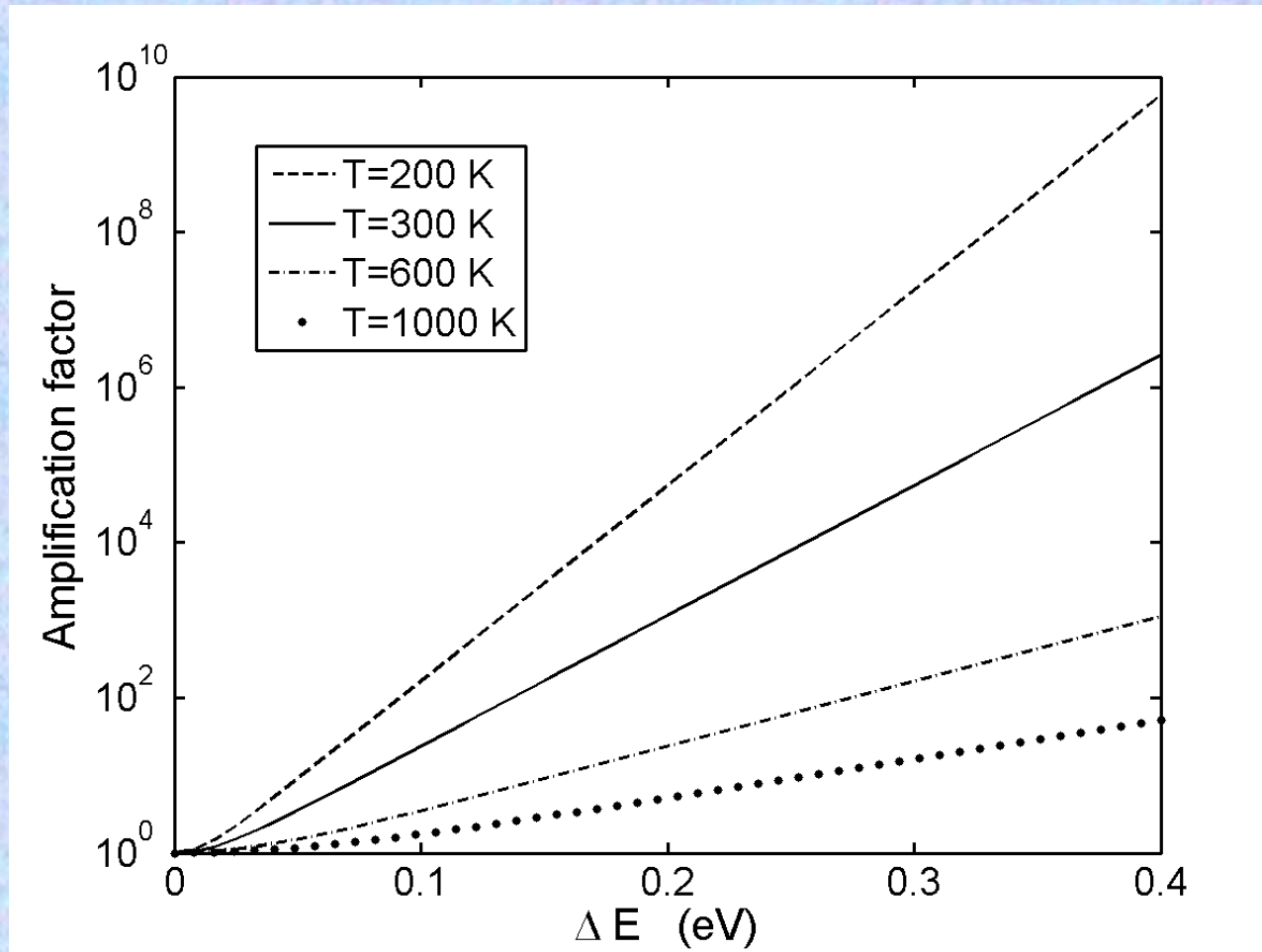
Suppose: E_A is changed in ΔE in Δt ; $-\Delta E$ in Δt

$$\left\langle \exp\left(-\frac{E}{k_B T}\right) \right\rangle = \frac{1}{2\Delta t} \left[\exp\left(-\frac{E_A + \Delta E}{k_B T}\right) \Delta t + \exp\left(-\frac{E_A - \Delta E}{k_B T}\right) \Delta t \right]$$
$$= \exp\left(-\frac{E_A}{k_B T}\right) \cosh\left(-\frac{\Delta E}{k_B T}\right) = \exp\left(-\frac{E_A}{k_B T}\right) I\left(-\frac{\Delta E}{k_B T}\right)$$

Amplification factor

$$I = \cosh\left(-\frac{\Delta E}{k_B T}\right) \approx \frac{1}{2} \exp\left(-\frac{\Delta E}{k_B T}\right)$$

ILMs: amplification effect on Arrhenius law

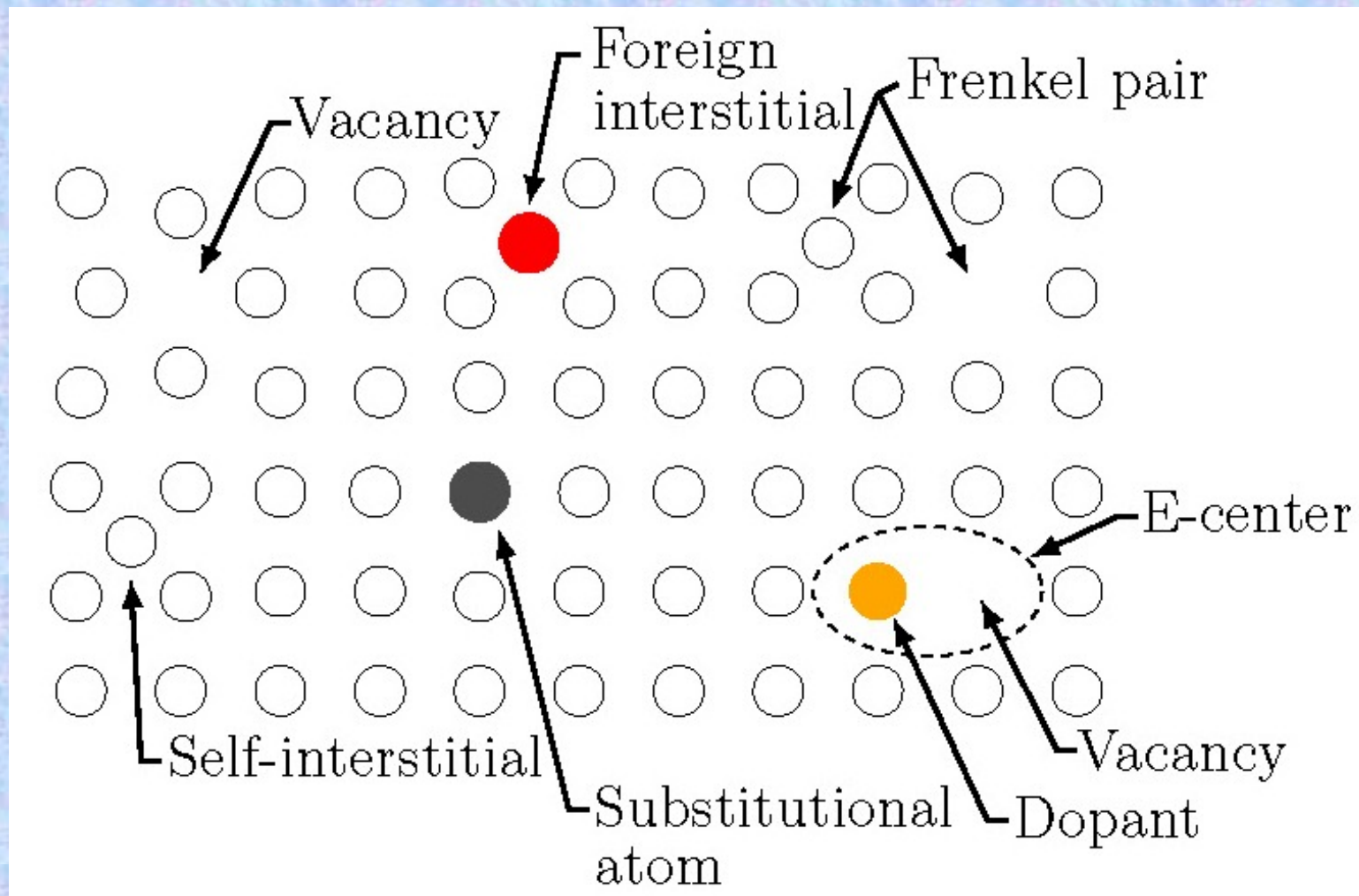


$$I = \cosh\left(-\frac{\Delta E}{k_B T}\right) \approx \frac{1}{2} \exp\left(-\frac{\Delta E}{k_B T}\right)$$

Dubinko, Selyshchev, Archilla,
Reaction-rate theory with account of the crystal anharmonicity.
Phys. Rev. E **83**, 041,124 (2011)

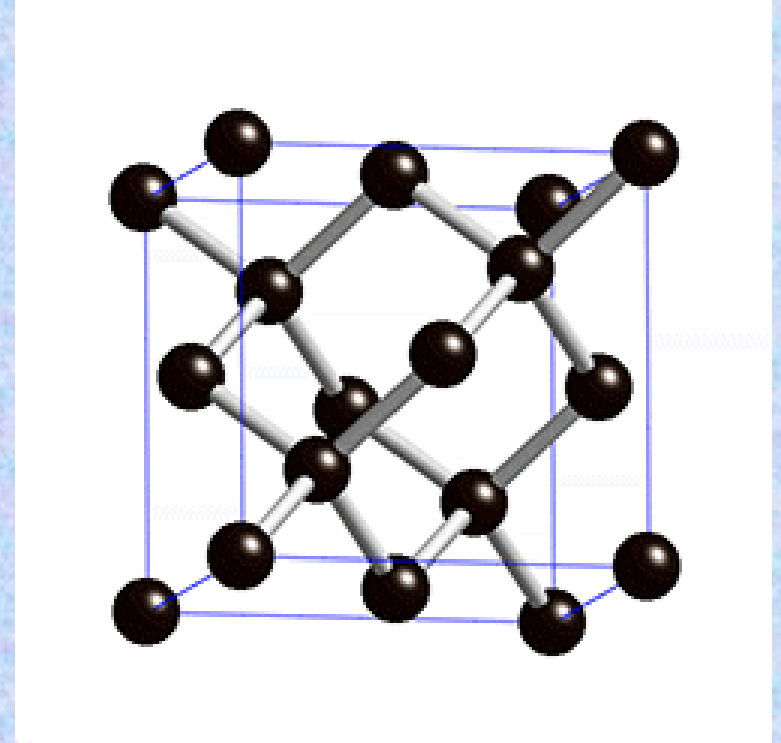
Defects in germanium

- Can be produced by irradiation
- Of technological interest

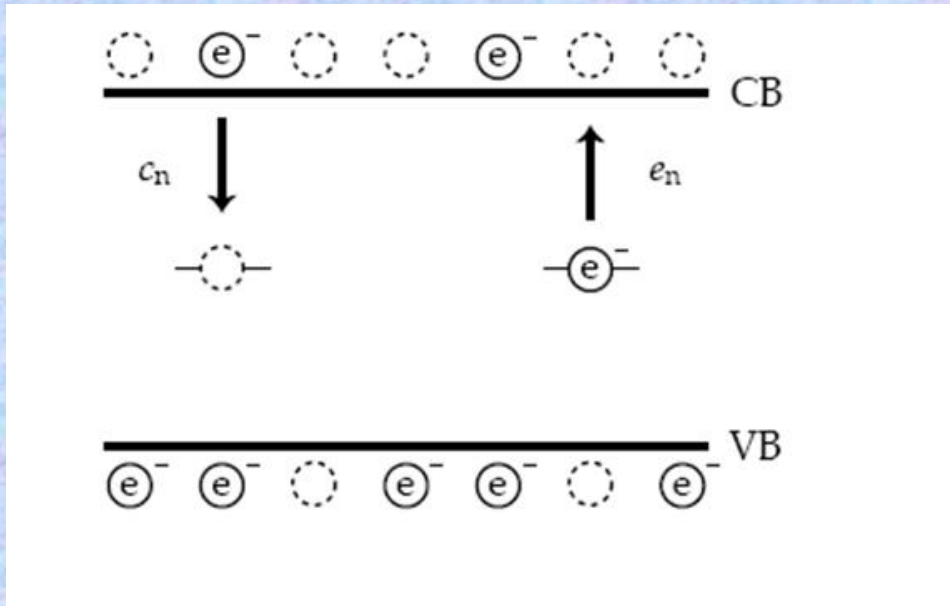


Some complex defects in germanium

- Di-vacancy V-V
- Tri-vacancy V_3 , Tetra-vacancy V_4
- Vacancy-Hydrogen VH_n
- I_2, I_3, \dots
- A center:
Foreign interstitial 0-Vacancy
- E center
Substitutional atom-Vacancy
Sb doped Ge: Sb-V



Defects as electron traps



Deep level electron trap:

$$E_T = E_t - E_c \geq 0.1 \text{ eV}$$

$$E_{0.37} : E_T = 0.37 \text{ eV}$$

$$c_n = \sigma_n \langle v_n \rangle n$$

Empty level. defect
can
-Capture an electron
from the CB

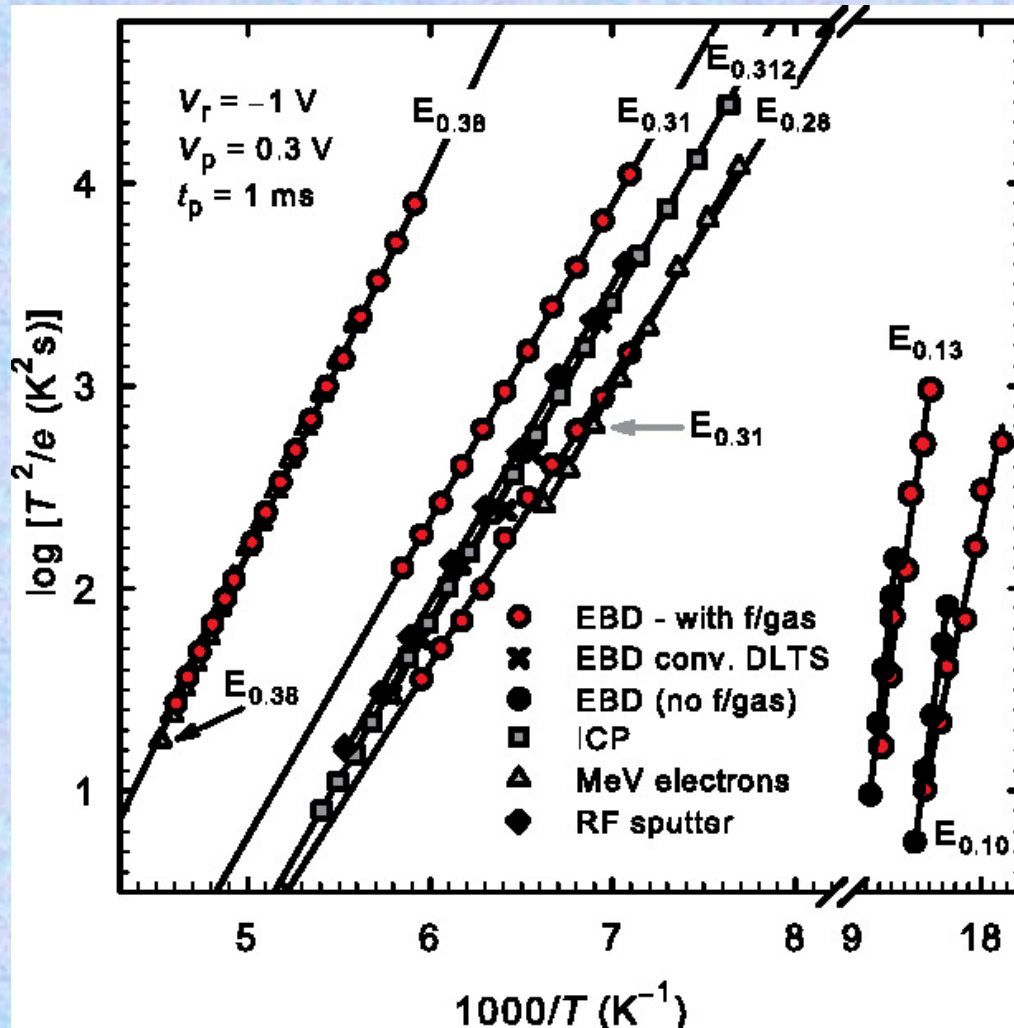
Defect populated by an
electron can :
-Emit an electron
to the CB

$$e_n = \sigma_n \langle v_n \rangle N_c \exp\left(-\frac{E_T}{kT}\right)$$

Capture rate c_n

Emission rate e_n

Signature of the trap



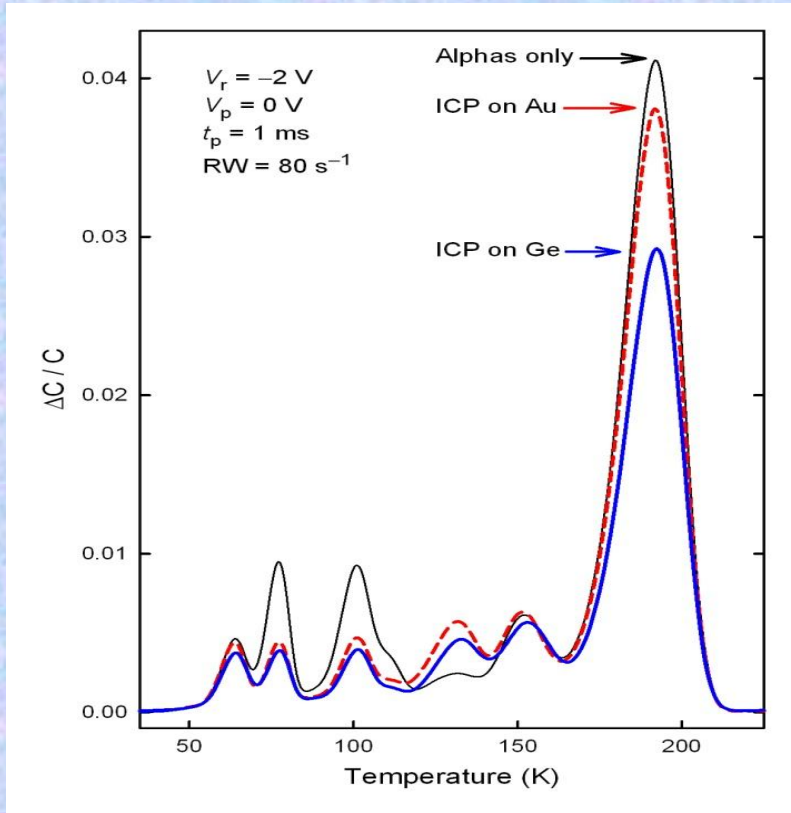
Energy of the trap

$$E_T$$

Capture cross section of the trap

$$\sigma_n$$

Plasma induced annealing and thermal annealing



$$\text{Ion-induced} \quad -\frac{dN_T}{dt} = K_i N_T$$

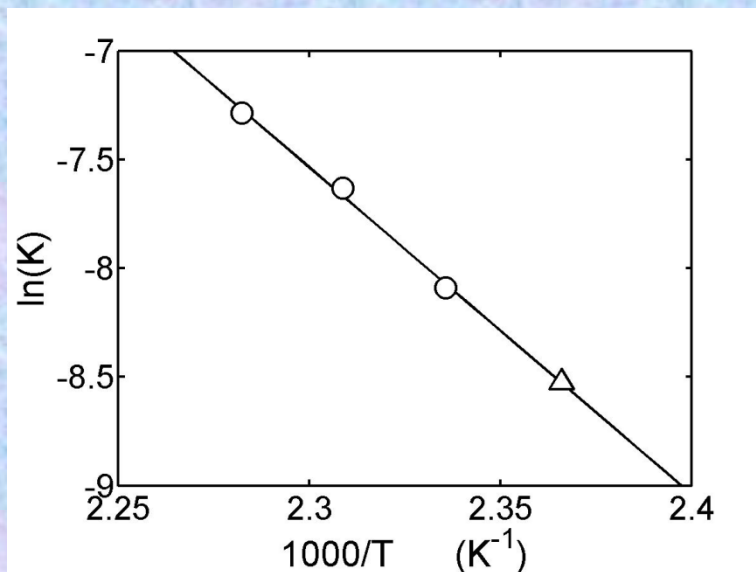
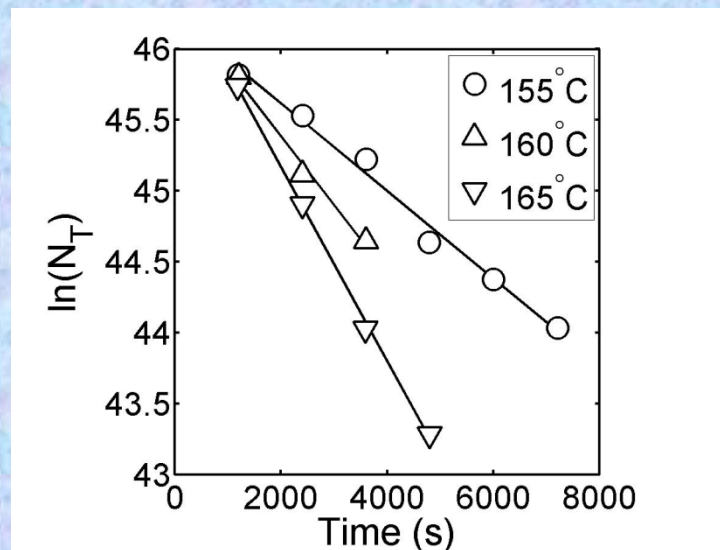
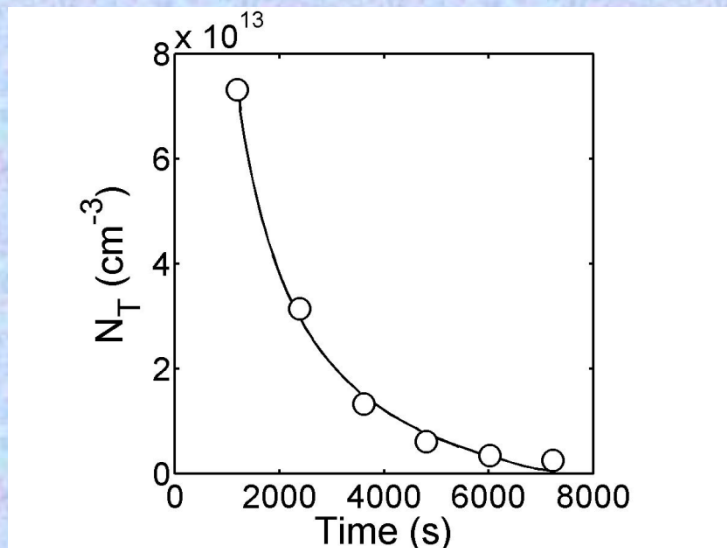
$$\text{Thermal} \quad -\frac{dN_T}{dt} = K_{th}(T) N_T$$

$$K_i = 2.0 \times 10^{-4} \text{ s}^{-1} \quad (\text{plasma})$$

$$K_{th} = 9.6 \times 10^{-11} \text{ s}^{-1} \quad (\text{thermal})$$

$$\text{Equivalent temperature} \quad T_A = 423 \text{ K} \quad (150 \text{ C}) \quad K_{th} = K_i$$

Thermal annealing

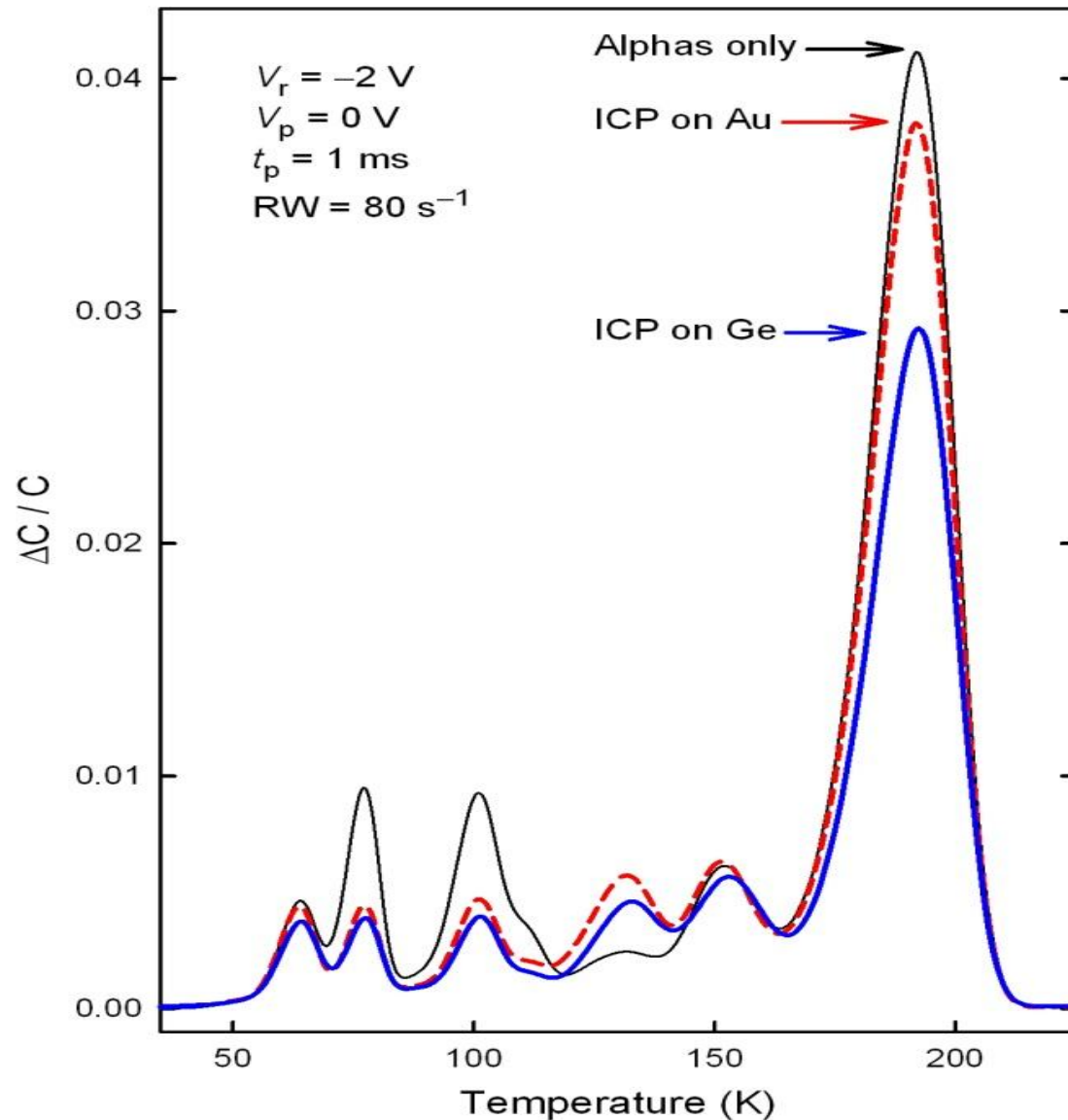


$$N_T = N_T(0)e^{-K_{th}t} \quad K_{th} = Ae^{-\frac{E_A}{k_B T}}$$

$$A = 0.55 \times 10^{12} \text{ s}^{-1}; \quad E_A = 1.3 \text{ eV}$$

$$T_A = 423 \text{ K} \quad (150 \text{ }^\circ\text{C})$$

Experimental details: 4 eV ICP plasma annealing



1.-Sb doped Ge is damage with 5 MeV alpha particles

2.- Rest 24 hours

3.-Au diode is evaporated in half the sample (half A)

4.- DLST in A
(black, alphas only)

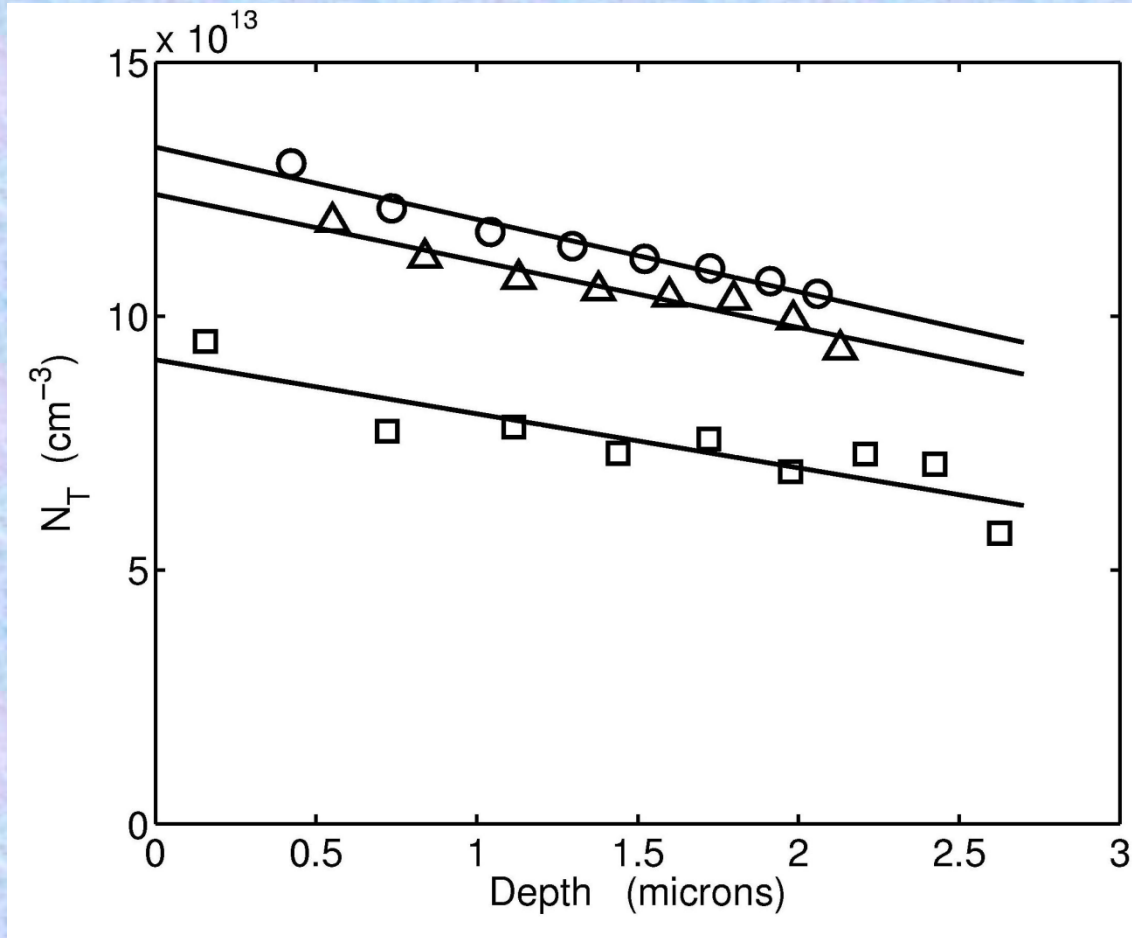
5.-ICP in A and B

6.- Au diode in B

7.- DLTS in A
(red, ICP on Au)

8.- DLTS in B
(blue, ICP on Ge)

Depth profile



Rate of annealing
decreases very
slowly

Our hypothesis: Ar ions impacting on Ge produce ILMs, which travel through Ge and anneal the defects.

Why?

- 1.-ILMs with MD in metals by Hyzhnyakov group have 0.5-5 eV
- 2.-The maximum energy transfer from Ar to Ge is 3.6 eV
- 3.- The activation energy for annealing an E center is 1.3 eV
- 4.- The energy to anneal a defect has to remain localized up to 10^4 lattice units
- 5.- Increasing the energy of the plasma does not enhance the effect, this is because ILMs typically have a definite range of energies.
- 6.- At least stationary ILMs have been obtained for Si and Ge with MD.

More is less: Some numbers on cross sections for ions

Ion current can be measured, then
ion cross section can be calculated

$$-\frac{dN_T}{dt} \approx \sigma_i \Phi_i N_T$$

For 4 eV Ar ions:

$$\Phi_i = 5.5 \times 10^{10} \text{ cm}^{-2} \text{ s}^{-1}$$

$$\sigma_i = 4.4 \sigma_0$$

for $\sigma_0 = (n_{Ge})^{-2/3} \approx 8 \text{ \AA}^2$

For 8 eV Ar ions:

$$\Phi_i(8 \text{ eV}) = 1.3 \times 10^{13} \text{ cm}^{-2} \text{ s}^{-1}$$

$$\sigma_i(8 \text{ eV}) = 0.0333 \sigma_0$$

$$\sigma_i(8 \text{ eV}) \approx \frac{1}{140} \sigma_i(4 \text{ eV})$$

Some numbers on energy of phonons and ILMs

1.-Ion current can be measured $-\frac{dN_T}{dt} \approx \sigma_i \Phi_i N_T$; $-\frac{dN_T}{dt} \approx \sigma_{ILM} \Phi_{ILM} N_T$

$$\sigma_i \Phi_i = \sigma_{ILM} \Phi_{ILM}$$

$$\Phi_{ILM} < \Phi_i \quad ; \quad \sigma_{ILM} > \sigma_i$$

$$n_{ILM} = \frac{\Phi_{ILM}}{c_s} \approx 10^{-18} n_{Ge} \quad E_{ILM} < E_{Ge,Max} \approx 3.6 \text{ eV}$$

Energy in phonons and energy in ILMs for same annealing

$$u_{ph} = \frac{1}{n_{Ge}} \int_{T_R}^{T_A} g(E) \langle n(E) \rangle E dE \approx 30.1 \text{ meV/atom}$$

$$u_{ILM} = \frac{1}{n_{Ge}} n_{ILM} E_{ILM} \approx 10^{-15} \text{ meV/atom}$$

$$\frac{u_{ILM}}{u_{ph}} \approx 10^{-16}$$

Conclusions:

0.- Plasma of 4eV produces annealing of defects very deep in Ge

Likely conclusions:

1.- 4 eV Ar hits produce ILM in Ge with very high efficiency

2.-ILMs of energy $\sim 3\text{eV}$ travel distances of the order of at least 10^4 lattice units

3.- The annealing efficiency of DB with respect to phonons is extremely large

4.- The energy delivered by a DB to a defect is ~ 0.3 eV

References

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Quodons in mica: nonlinear localized travelling excitations in crystals.

J.F.R Archilla *et al* (eds). Springer 2015 (to appear)

Chapters:

4 in mica; Russell, Eilbeck, Archilla, Russell

4 in MD of breathers (2D): Malomed, Cisneros, Johansson, Wattis

3 in MD of breathers (3D): Dmitriev, Hizhnyakov, Kosevich

3 in electrons and lattices: Brizhik, Chetverikov, Ebeling, Cruzeiro, Velarde

3 in semiconductors: Coelho, Dubinko, Archilla

3 in rogue waves, proteins, BC condensates

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Experimental observation of intrinsic localized modes in germanium.-

-Coelho, S.M.M., Archilla, J.F.R., Auret, F.D., Nel, J.M.: **The origin of defects induced in ultrapure germanium by electron beam deposition.**

-Dubinko, V.I., Archilla, J.F.R., Dmitriev, S.V., Hizhnyakov, V.: **Rate theory of acceleration of the defect annealing driven by discrete breathers.**

-Hizhnyakov, V., Haas, M., Shelkan, A., Klopov, M.: **Standing and moving discrete breathers with frequencies above the phonon spectrum.**

NONLINEAL 2016: Sevilla, June 7-10, 2016

Conference on nonlinear mathematics and physics

