



Bifurcation Analysis of Discrete Breathers in a Nonlinear bent chain of oscillators

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Outline

- Concept of discrete breather.
- Discrete breathers in homogeneous lattices.
Bifurcations.
- Discrete breathers in bent chains.
Bifurcations.
- Conclusions.

Discrete breathers

MOVIE

Discrete breathers

- Different frameworks:

- Discrete Nonlinear Klein–Gordon Equation:

$$\ddot{u}_n + V'(u_n) + C \sum_m W'(u_n, u_m) = 0$$

- Fermi–Pasta–Ulam model: $V(u_n) = 0 \forall n$.

- Discrete Nonlinear Schrödinger Equation:

$$i\dot{u}_n + \gamma|u_n|^2 u_n - C(2u_n - u_{n+1} - u_{n-1}) = 0$$

- $u_n \in \mathcal{E}_s^2(\omega_b)$; $V, W \in \mathcal{C}^2$; $n \in \mathbb{Z}^d$.

MacKay-Aubry theorem

- R.S. MacKay and S. Aubry. *Nonlinearity* 7 (1994) 1623-1643.
- *If a periodic orbit of $H = p^2/2 + V(u)$ with action I_0 is non-resonant and anharmonic, then the periodic orbit of $\ddot{u}_n + V'(u_n) + C(2u_n - u_{n+1} - u_{n-1}) = 0$ at $C = 0$ given by $x_0(t) = X(I_0, \omega_b(I_0)t)$ and $x_n(t) = 0$ for all sites $n \neq 0$ has a locally unique continuation as a periodic orbit of $\ddot{u}_n + V'(u_n) + C(2u_n - u_{n+1} - u_{n-1}) = 0$ of the same period $T = 2\pi/\omega_b(I_0)$ for C small enough.*
- *The C^1 norm of the oscillation on site n for that periodic solution decays exponentially as $n \rightarrow \infty$, as long as DF remains invertible.*

Linear stability

- Dynamical equations of a perturbation $u_n \rightarrow u_n + \xi_n$ ($\xi \in \mathcal{C}^2$):

$$\ddot{\xi}_n(t) + V''(u_n(t))\xi_n(t) - C \sum_i W''(u_{n+i}(t) - u_n(t))(\xi_{n+i}(t) - \xi_n(t)) + \\ + C \sum_i W''(u_n(t) - u_{n-i}(t))(\xi_n(t) - \xi_{n-i}(t)) = 0$$

- Linear stability \rightarrow Floquet analysis:

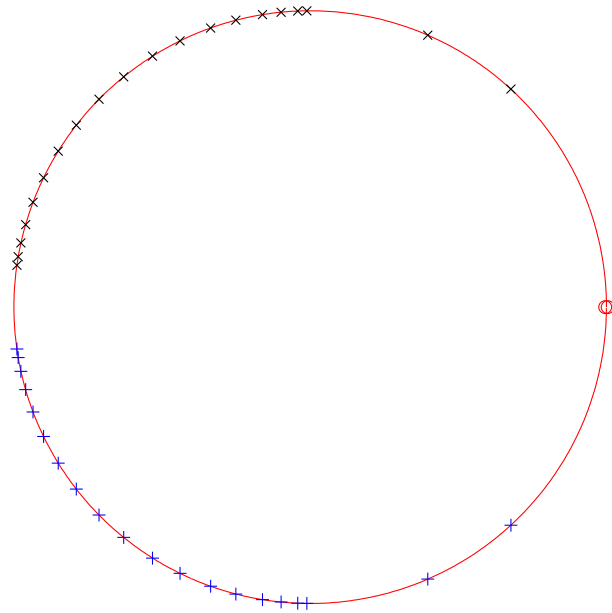
$$\begin{bmatrix} \xi(T) \\ \dot{\xi}(T) \end{bmatrix} = \mathcal{F}_o \begin{bmatrix} \xi(0) \\ \dot{\xi}(0) \end{bmatrix}$$

- Floquet's operator is real and symplectic:

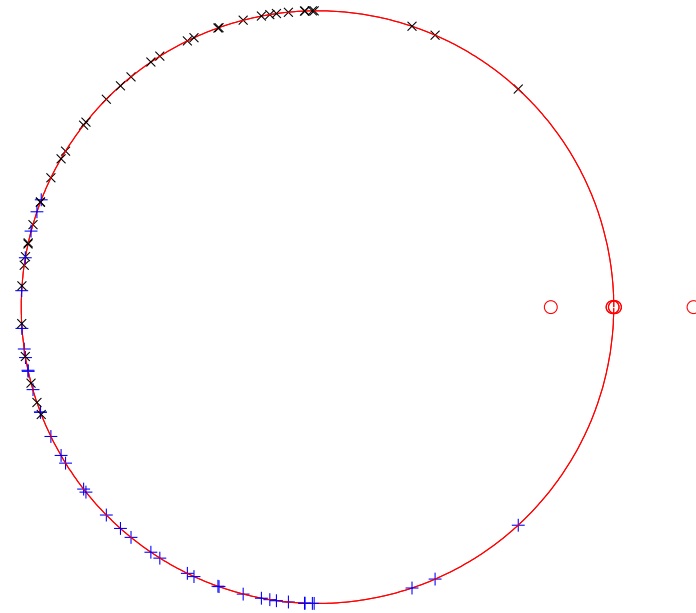
- $\lambda, 1/\lambda, \lambda^*, 1/\lambda^*$ are eigenvalues.

- A discrete breather in a Hamiltonian Klein–Gordon lattice is linear stable if and only if all eigenvalues are on the unit circle.

Linear stability

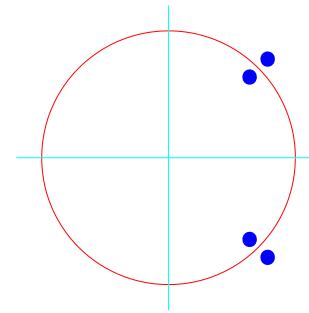
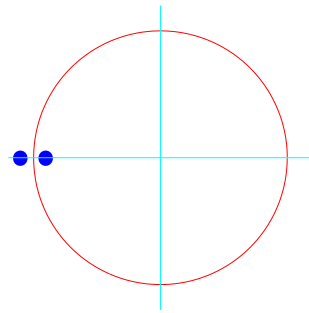
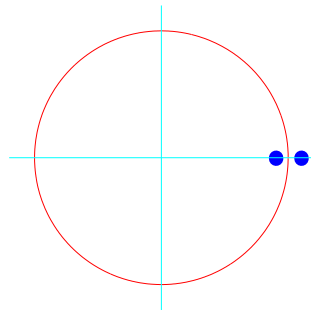
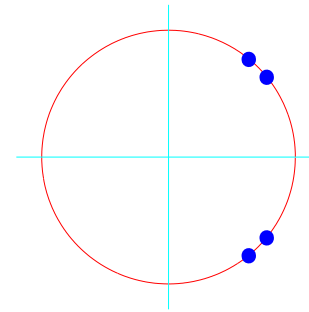
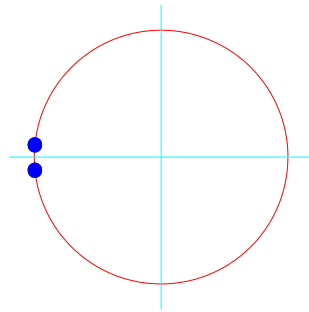
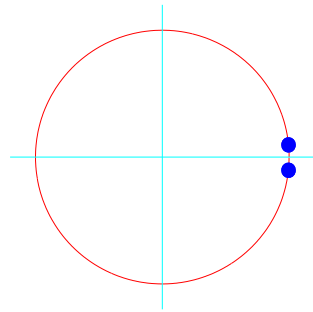


Stable



Unstable

Bifurcations

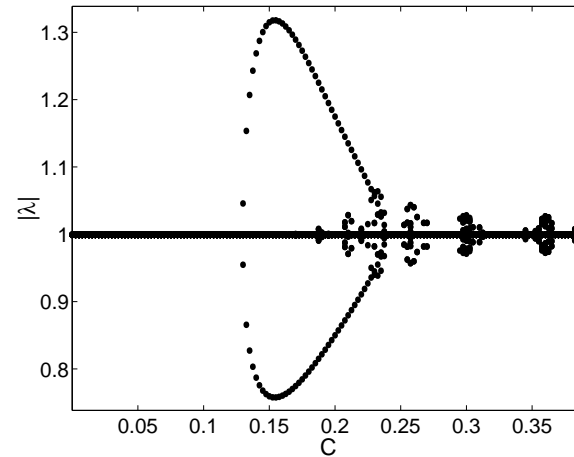
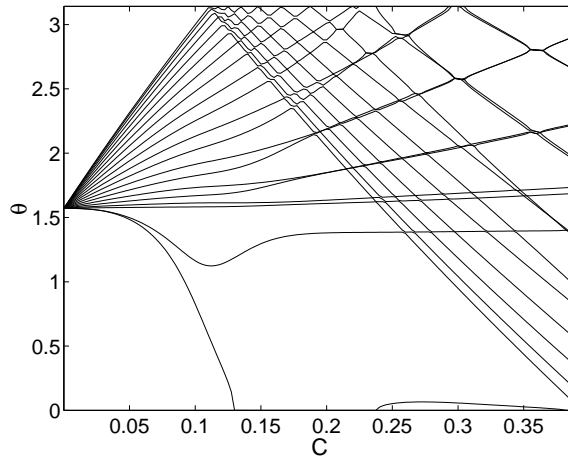


Harmonic

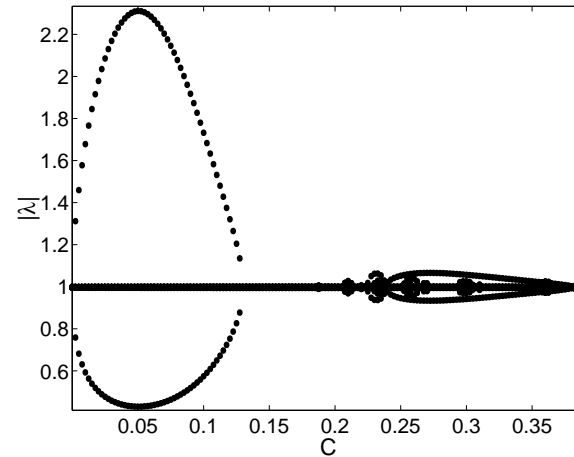
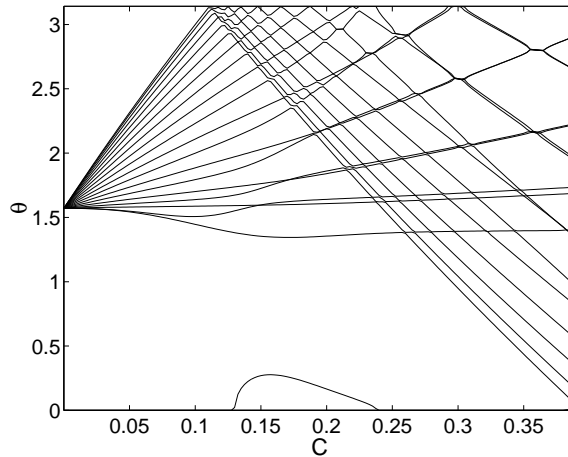
Subharmonic

Oscillatory

Stability exchange (Marginal mode)



Site-centered breather



Bond-centered breather

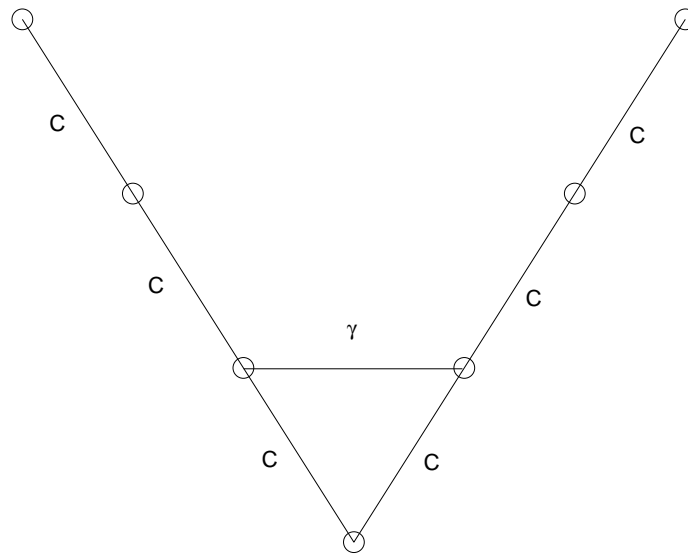
Mobile breather

MOVIE

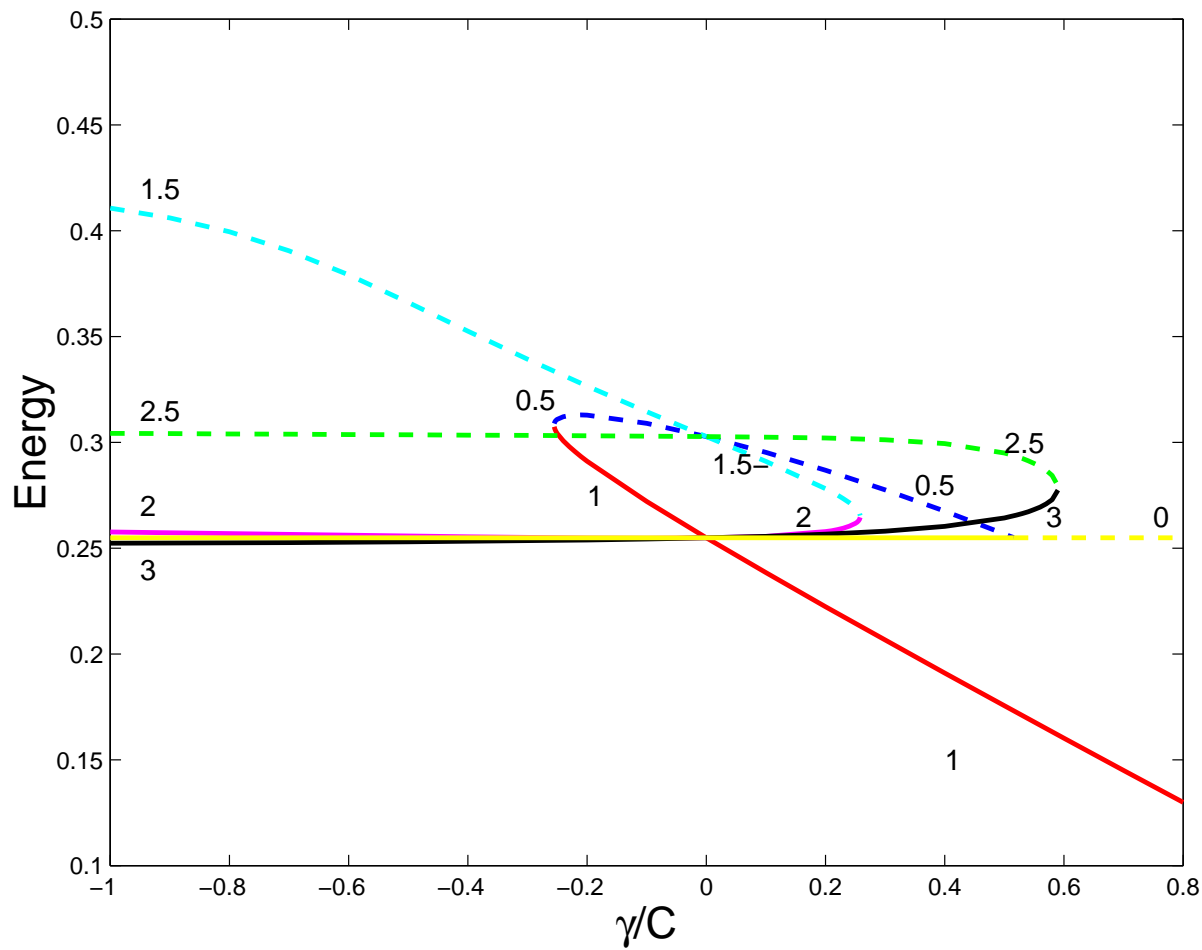
Bent (wedged) chain

$$\ddot{u}_n + V'(u_n) + C \sum_m (2u_n - u_{n+1} - u_{n-1}) + \gamma[(u_n - u_{n-2})\delta_{n,1} + (u_n - u_{n+2})\delta_{n,-1}] = 0$$

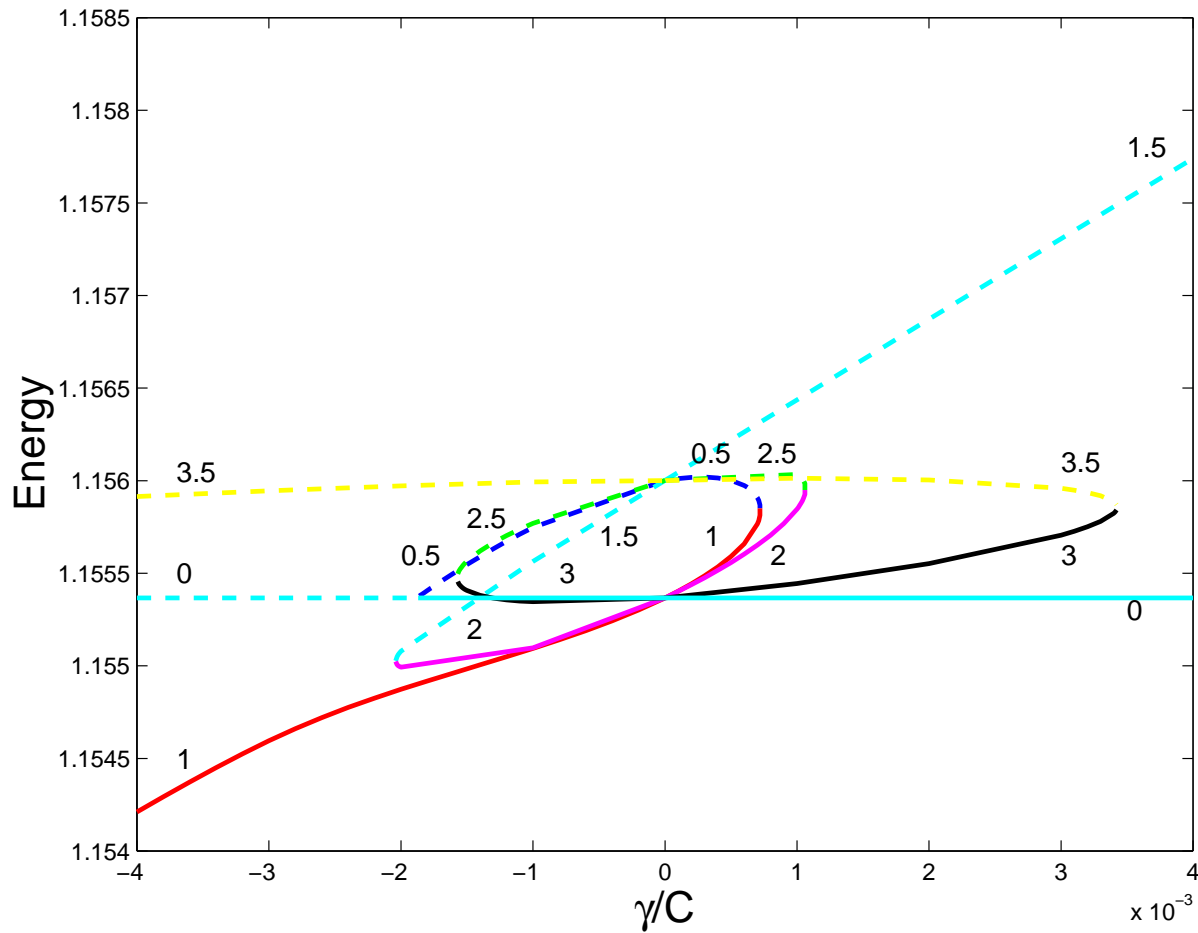
- u_n : Out-of-plane displacements.
- $n = 0$: Vertex of the wedge.



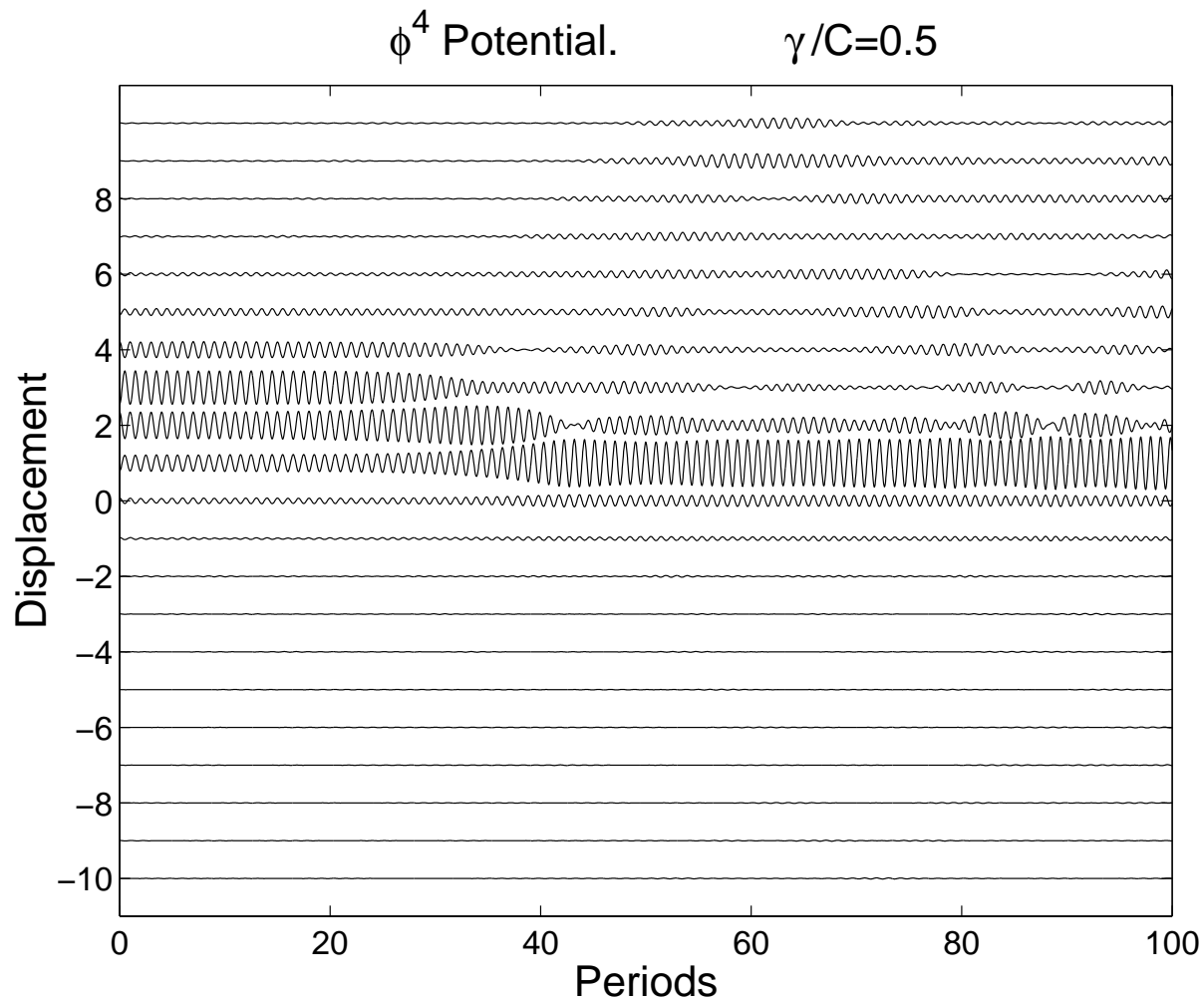
Bifurcations. ϕ^4 potential



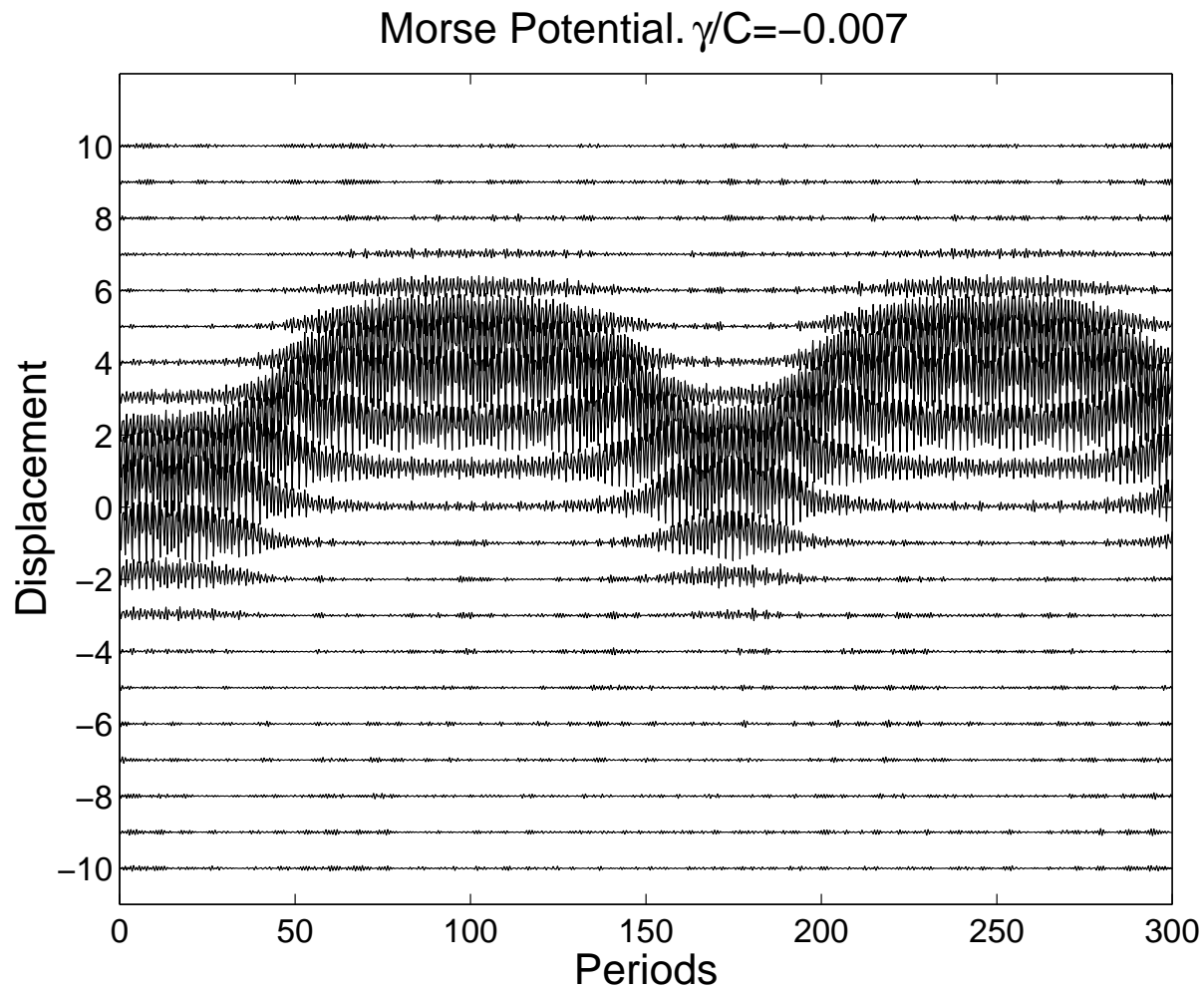
Bifurcations. Morse potential



Switching in ϕ^4 potential



Mobility in Morse potential



Conclusions

- Homogeneous Klein–Gordon lattices: Summary of (stability) bifurcations.
- Wedged Klein–Gordon chains: Ground state changes with γ .
 - ϕ^4 potential: Switching.
 - Morse potential: Spontaneous mobility.
- Nonlinear Physics Group website:

<http://www.us.es/gfnl>