

Nonlinear charge transport in DNA mediated by twist modes



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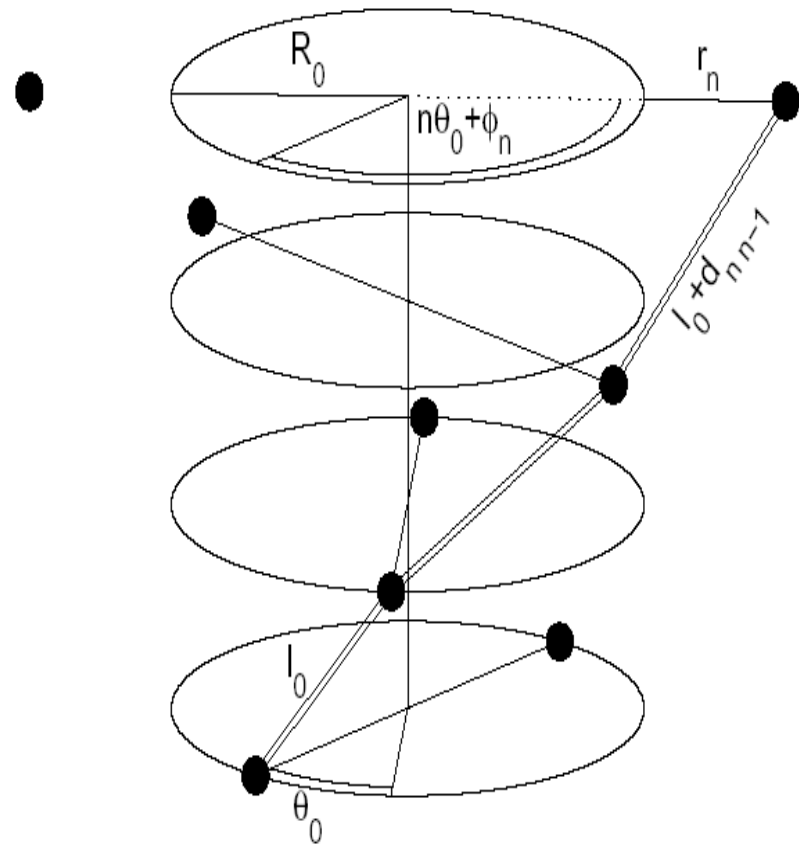
Charge transport in DNA



- ⌘ Important role for biological functions. Biosynthesis and DNA repair after radiation damage.
- ⌘ Technological interest. Possibility of building electronic devices based on biomaterials.

The model

- ⌘ Modification of the *twist-opening* model.
- ⌘ Bond potential treated as harmonic ones.
- ⌘ Three dimensional, semi-classical, tight-binding model.



Hamiltonian

- Lattice oscillators treated classically

$$H_{rad} = \sum_n \left[\frac{1}{2M} (p_n^r)^2 + \frac{M\Omega_r^2}{2} r_n^2 \right],$$

$$H_{twist} = \sum_n \left[\frac{1}{2J} (p_n^\phi)^2 + \frac{J\Omega_\phi^2}{2} (\phi_n - \phi_{n-1})^2 \right]$$

- Charge described by a tight-binding system

$$\widehat{H}_{el} = \sum_n E_n |n\rangle \langle n| - V_{n-1,n} |n-1\rangle \langle n| - V_{n+1,n} |n+1\rangle \langle n|$$

$$E_n = E_n^0 + kr_n \quad V_{n,n-1} = V_0(1 - \alpha d_{n,n-1})$$

$$d_{n,n-1} = [a^2 + (R_0 + r_n)^2 + (R_0 + r_{n-1})^2 - 2(R_0 + r_n)(R_0 + r_{n-1}) \cos(\theta_0 + \theta_{n,n-1})]^{1/2} - l_0$$

Dynamical equations

$$\begin{aligned}
 i\tau\dot{c}_n &= (E_n + k r_n) c_n \\
 &\quad - (1 - \alpha d_{n+1,n}) c_{n+1} - (1 - \alpha d_{n,n-1}) c_{n-1}, \\
 \ddot{r}_n &= -r_n - k |c_n|^2 \\
 &\quad - \alpha \left[\frac{\partial d_n}{\partial r_n} (c_n^* c_{n-1} + c_n c_{n-1}^*) + \frac{\partial d_{n+1}}{\partial r_n} (c_{n+1}^* c_n + c_{n+1} c_n^*) \right], \\
 \ddot{\phi}_n &= -\Omega^2 (2\phi_n - \phi_{n-1} - \phi_{n+1}) \\
 &\quad - \alpha V \left[\frac{\partial d_n}{\partial \phi_n} (c_n^* c_{n-1} + c_n c_{n-1}^*) + \frac{\partial d_{n+1}}{\partial \phi_n} (c_{n+1}^* c_n + c_{n+1} c_n^*) \right]
 \end{aligned}$$

- Scaled parameters

$$\tau = 0.053, \Omega^2 = 0.013, V = 2.5 \times 10^{-4}$$

$$R_0 = 63.1, l_0 = 44.5 \quad \text{No reliable data} \quad k=1, \alpha \text{ adjustable}$$

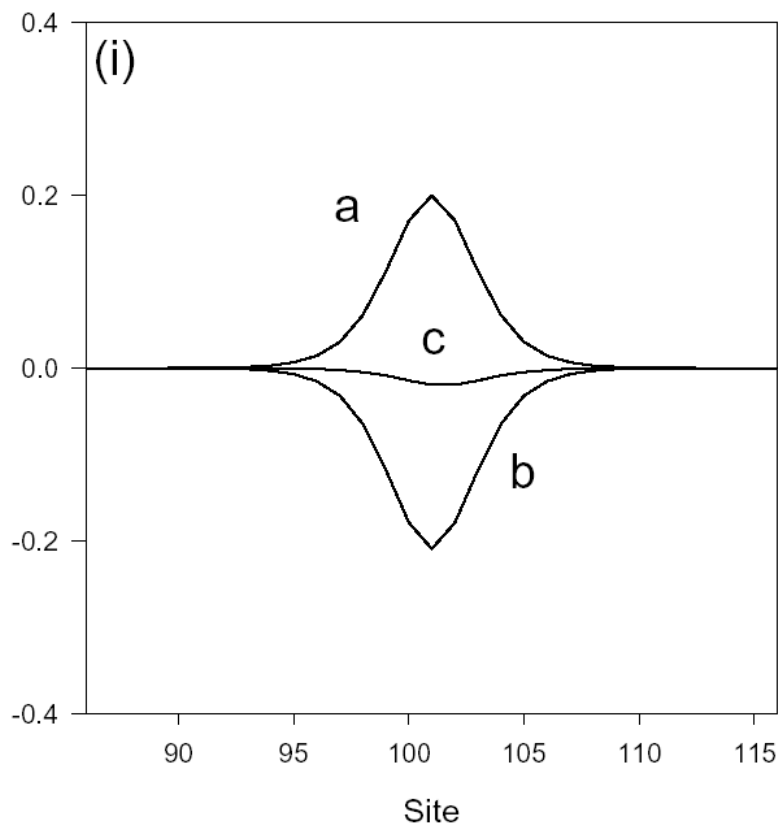
Stationary polaron-like states

- τ small.
- Born-Oppenheimer approximation.
- Stationary solutions must be attractors of the map:

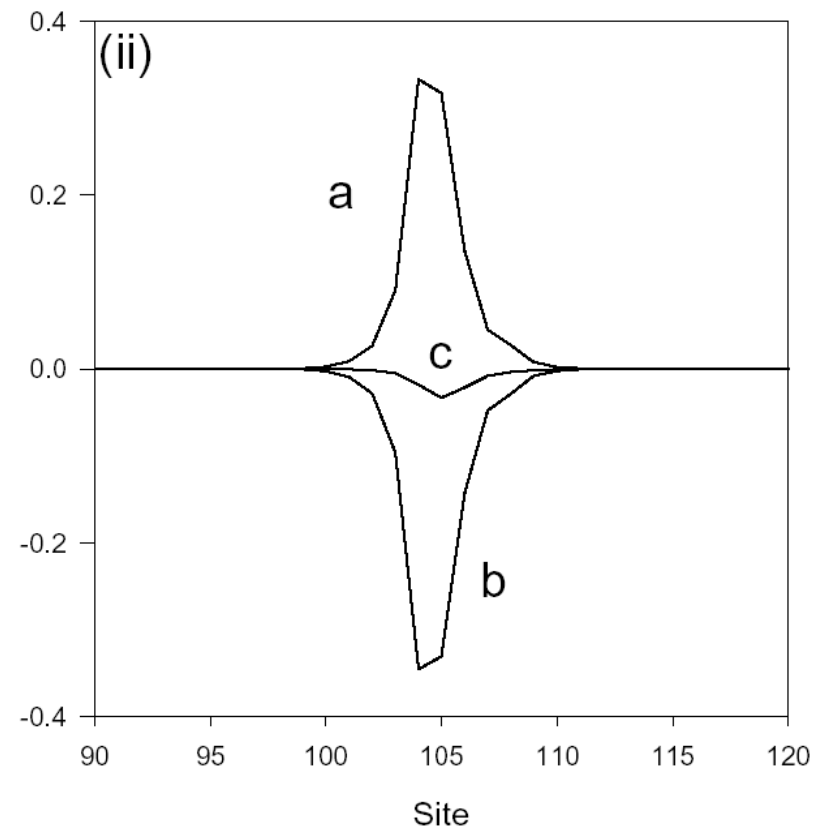
$$\begin{aligned}
 r'_n &= -k |c_n|^2 \\
 &\quad - \alpha \left[\frac{\partial d_{n,n-1}}{\partial r_n} (c_n^* c_{n-1} + c_n c_{n-1}^*) + \frac{\partial d_{n+1,n}}{\partial r_n} (c_{n+1}^* c_n + c_{n+1} c_n^*) \right], \\
 \phi'_n &= \frac{1}{2} (\phi_{n+1} + \phi_{n-1}) \\
 &\quad - \frac{\alpha V}{2\Omega^2} \left[\frac{\partial d_{n,n-1}}{\partial \phi_n} (c_n^* c_{n-1} + c_n c_{n-1}^*) + \frac{\partial d_{n+1,n}}{\partial \phi_n} (c_{n+1}^* c_n + c_{n+1} c_n^*) \right], \\
 c'_n &= \frac{[(E_n + k r'_n) c_n - (1 - \alpha d'_{n+1,n}) c_{n+1} - (1 - \alpha d'_{n,n-1}) c_{n-1}]}{\|(E_n + k r'_n) c_n - (1 - \alpha d'_{n+1,n}) c_{n+1} - (1 - \alpha d'_{n,n-1}) c_{n-1}\|}, \\
 d' &= d(r', \phi')
 \end{aligned}$$

Stationary states

⌘ Ordered case $E_n^0 = E_0$

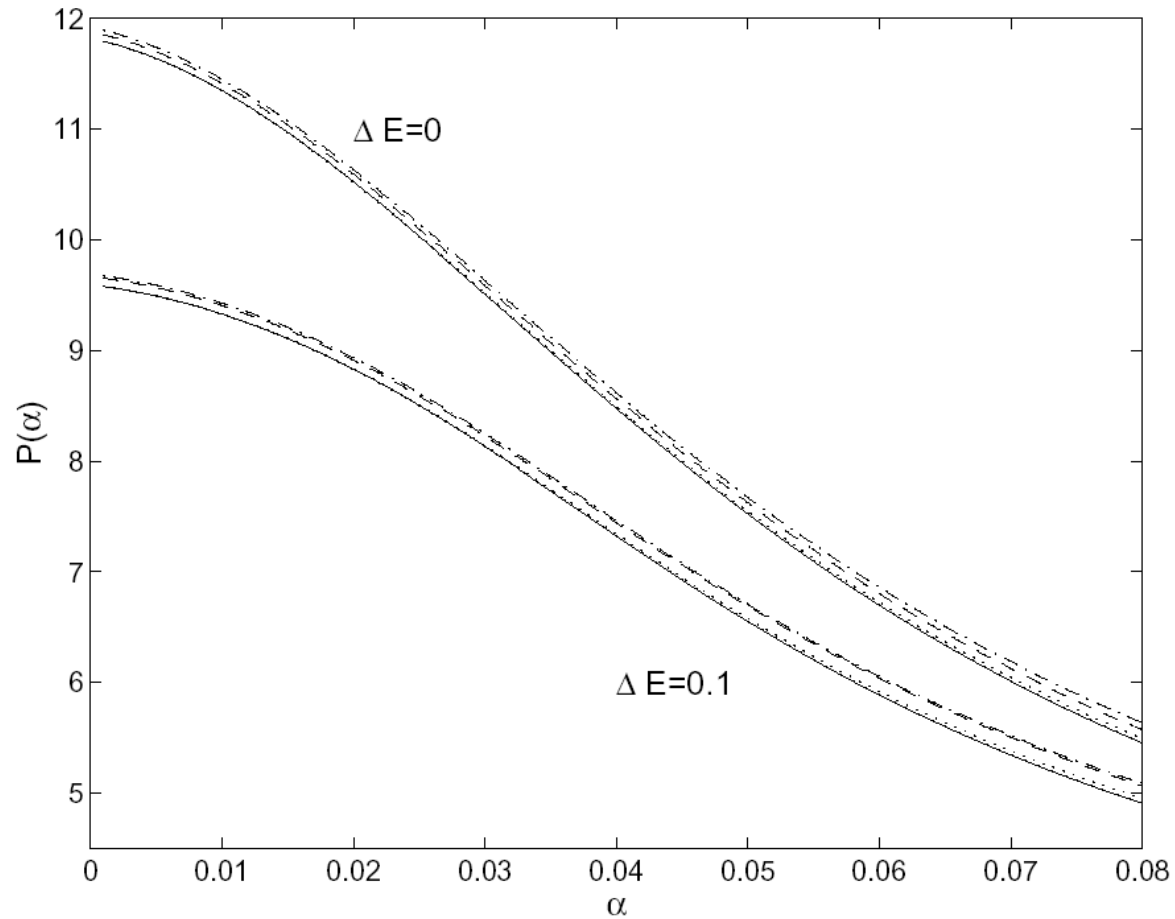


⌘ Disordered case $E_n^0 \in [-\Delta E_0, \Delta E_0]$




a) Electronic amplitude; b) Radial displacements, c) Angular twist

Participation number

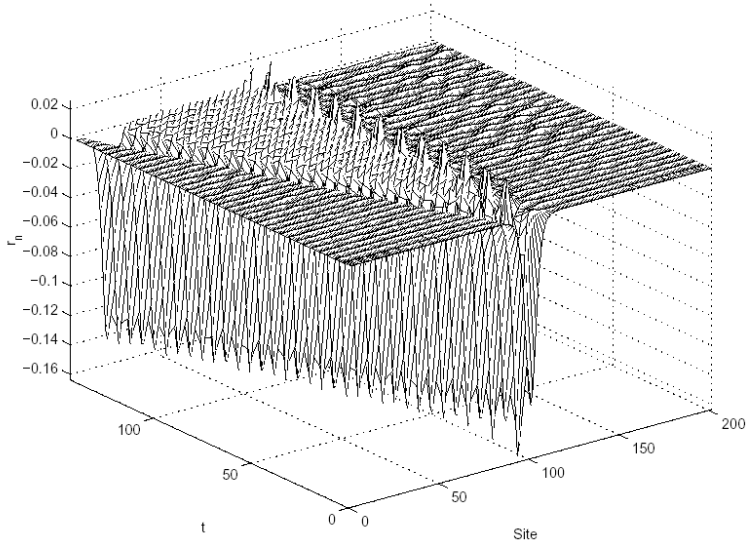
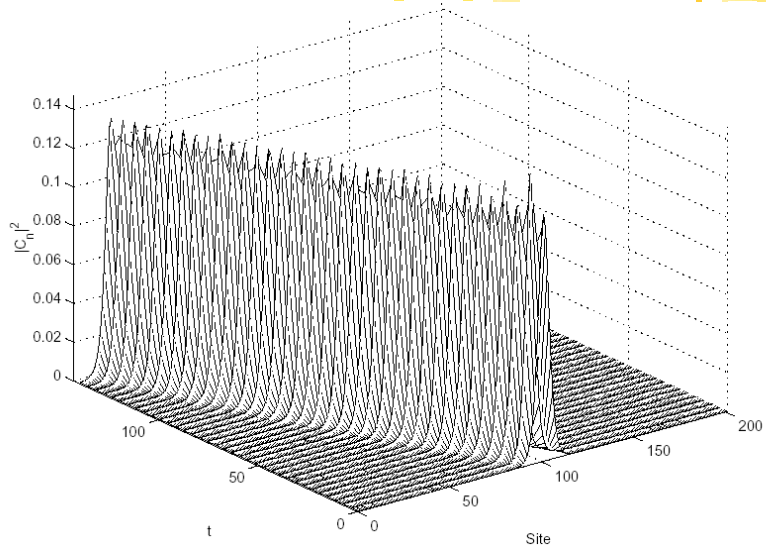


Charge transport in the absence of disorder

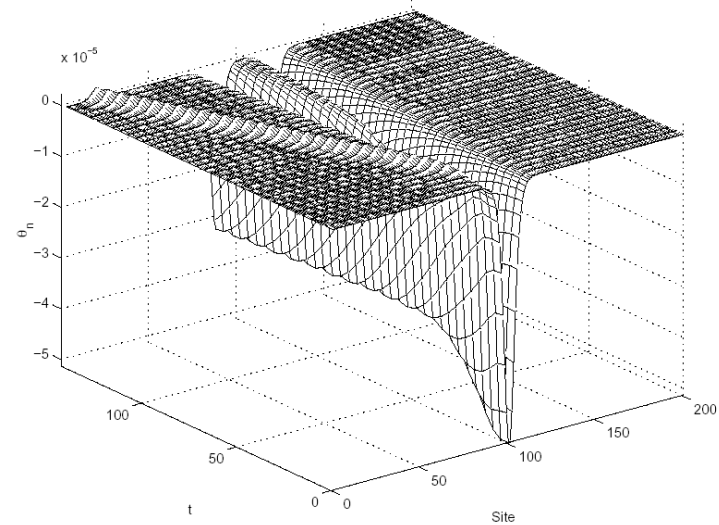


- Stationary solutions can be moved (under certain conditions).
- Discrete gradient method.
- Three different regimes depending on parameter α
 - Radial movability regime
 - Twist movability regime
 - Mixed regime

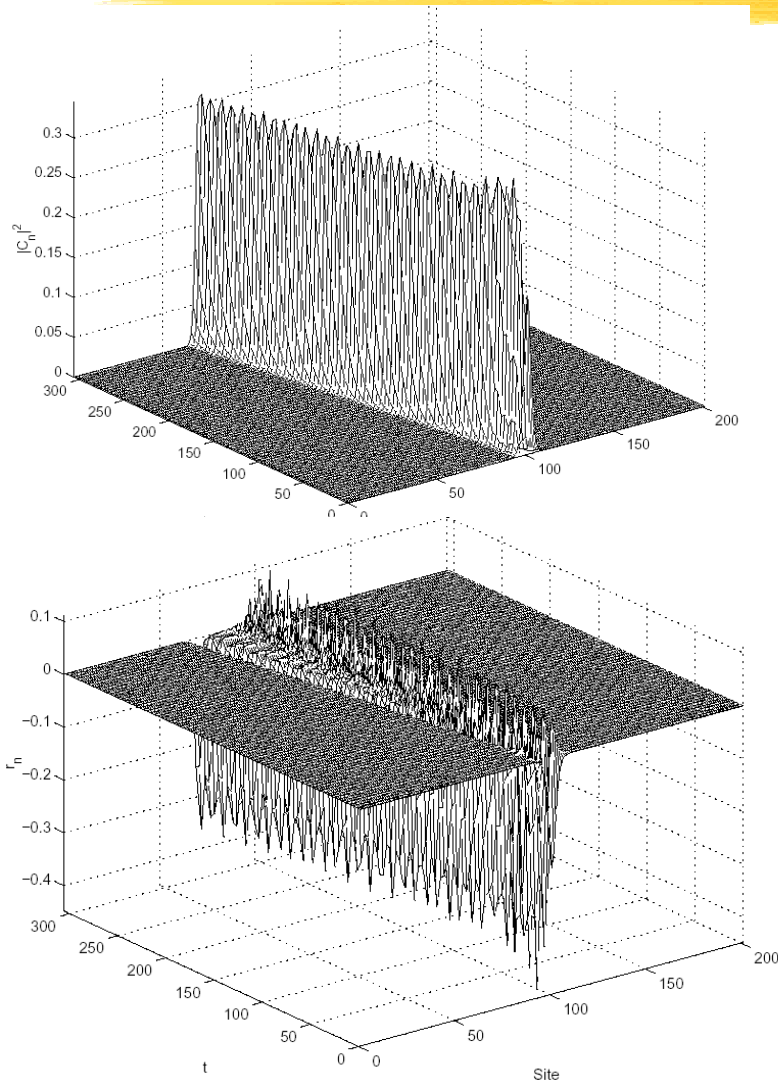
Ordered case I. Radial regime



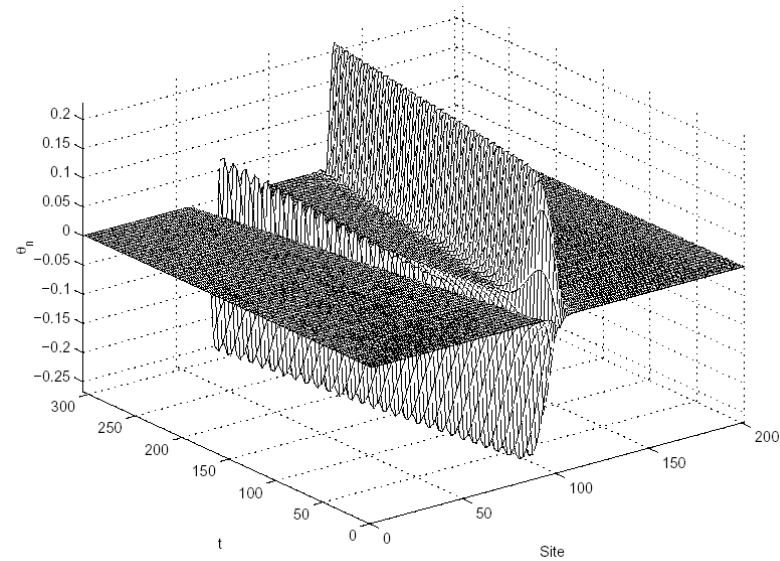
- α small enough (~ 0.0002).
- Velocity increases with perturbation.



Ordered case II. Twist regime



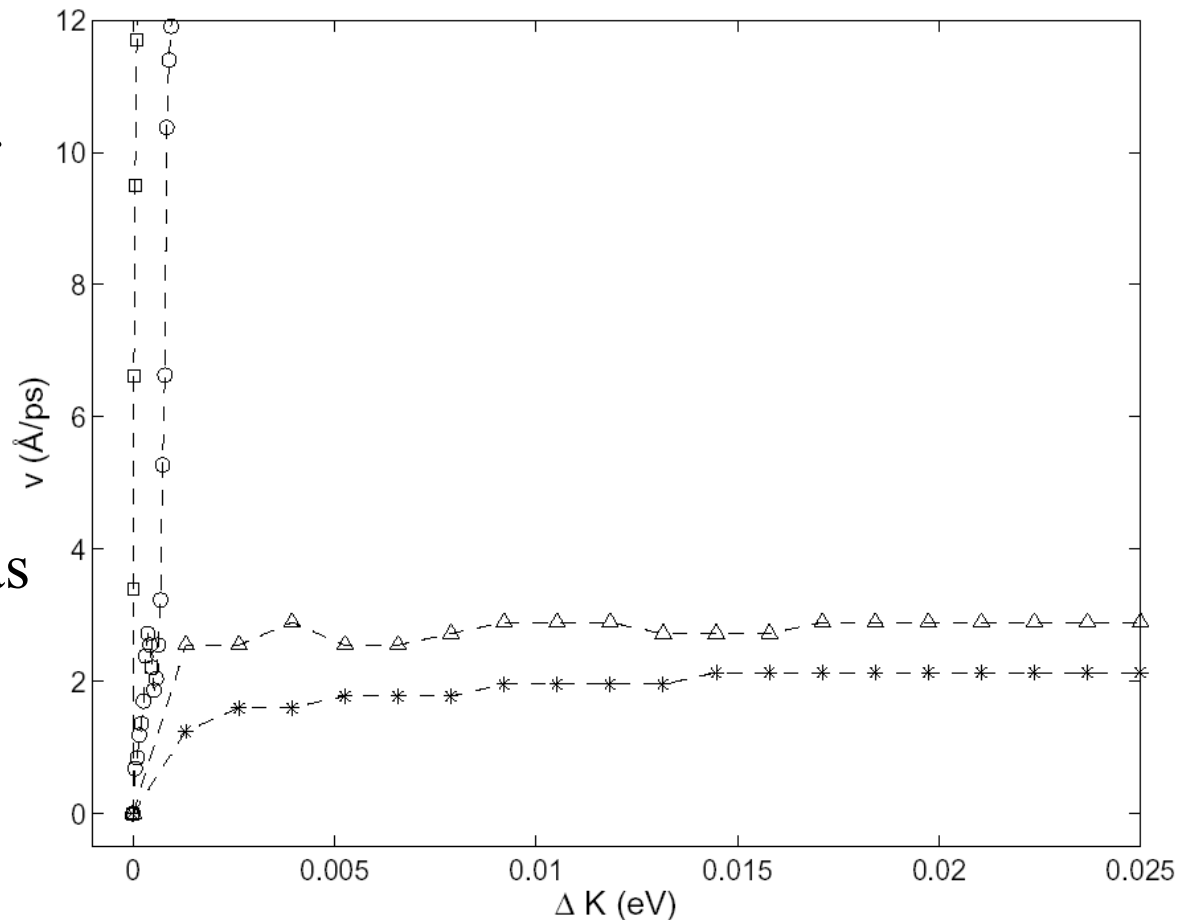
- Larger values of parameter α (~ 0.05).
- There exist a limit velocity.



Ordered case III. Mixed regime

Mixed regime for intermediate values of parameter α .

- Radial movability requires less energy.
- Radial movability has higher velocity.

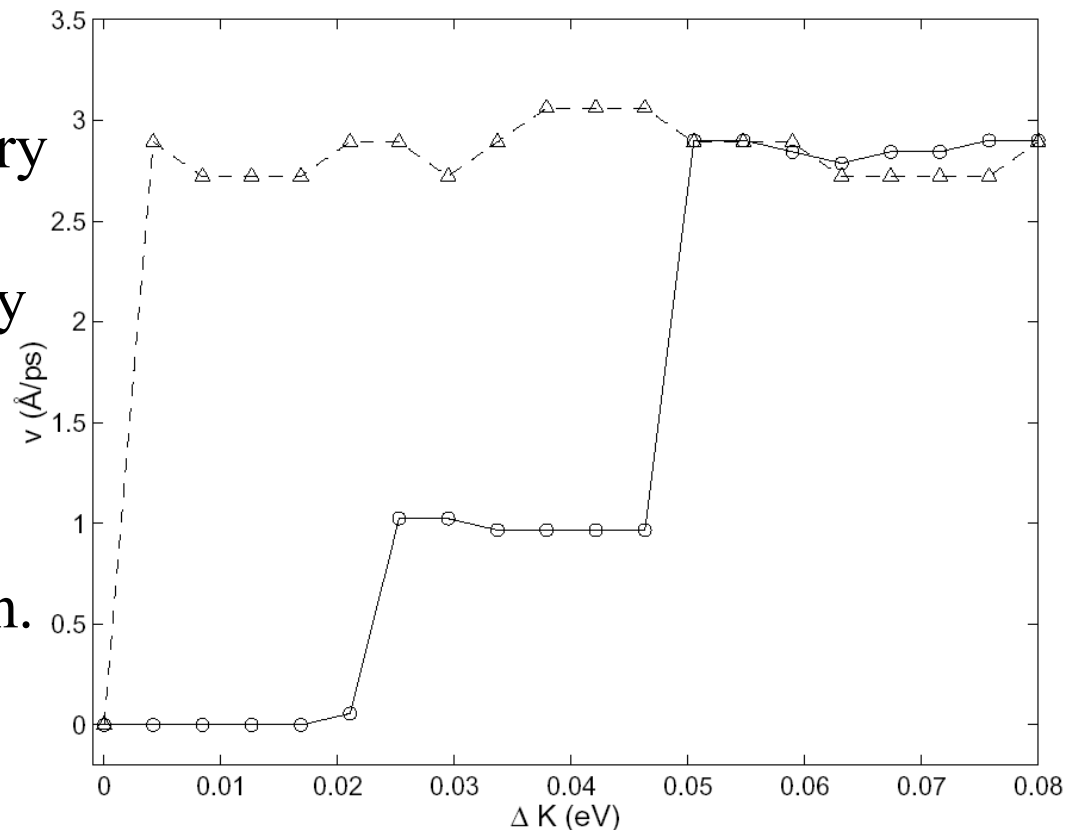


Squares: Radial regime; Circles: Mixed regime, radial activation;
Triangles: Mixed regime, angular activation; Stars: Twist regime

Palmero et al, Les Houches 2003

Disordered case

- Radial movability regime very sensitive to disorder.
- Twist movability regime very robust to disorder.
- Mixed regime, angular activation leads more robust polarons than radial activation.



Circles: Disordered case; Triangles:
Ordered one

Conclusions



- Analysis of the influence of the radial and angular perturbations on the properties of moving polarons in a three-dimensional, semi-classical, tight-binding model for DNA in both the ordered and disorderd case.
 - Three regimes in function of the coupling of the transfer integral with the deformations of the hydrogen bonds.
 - Properties of the moving polaron different in each regime.
- Mobility induced by angular activation is more robust with respect to parametric disorder, has lower velocity and the activation energy is higher than in radial movability regime.

Further works



- ⌘ Analysis of the influence of an impurity in the movement of the polaron.
- ⌘ We consider a different well depth (in radial variables) in a site.
 - ☒ Radial activation: Polaron is reflected or trapped by the impurity.
 - ☒ Twist activation: Polaron cross the impurity.

References



- ⌘ M. Barbi. Localized Solutions in a model of DNA Helicoidal Structure. PhD Thesis, Universitàt degli Studi di Firenze, 1998.- M. Barbi, S. Cocco and M. Peyrard. Phys. Lett. A, 253(5-6):358,1999.- M. Barbi, S. Cocco, M. Peyrard, and S. Ruffo. Jou. Biol. Phys.,24:97, 1999.
- ⌘ D. Hennig, J.F.R. Archilla, and J. Agarwal. Physica D, in press 2003.- J.F.R. Archilla, D.Hennig, and J. Agarwall. To appear. Nlin. PS/0211026, 2003.
- ⌘ F. Palmero, J.F.R. Archilla, D. Henning, and F.R. Romero. Preprint.