Creation and annealing of point defects in germanium crystal lattices by subthreshold

energy events

ergio M. M. Coelho¹, Juan F. R. Archilla² and F. Danie Auret ¹Physics Department, University of Pretoria, South Africa ² Group of Nonlinear Physics, University of Sevilla, Spain



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Our team in Pretoria





Juan FR Archilla (Project leader) Group of Nonlinear Physics Lineal (GFNL) University of Sevilla, Spain With Sergio Coelho, F Danie Auret Department of Physics University of Pretoria, South Africa Vladimir Dubinko Kharkov Physical-Technical Institute, Kharkov, Ukraine Vladimir Hizhnyakov Institute of Physics, University of Tartu Tartu, Estonia













sergio@up.ac.za



Outline

- Why study this? Low energy? Really?
- Germanium ultrapure material
- Hydrogen in germanium
- DLTS Deep Level Transient Spectroscopy
- An experiment: ICP annealing of E-center

Summary and conclusions

Why low energy?



Why study this? Low energy? Really?

- Impacts generate Moving Intrinsic Localized Modes (ILMs) or Discrete Breathers (DBs) in Ge
- ILMs in Ge Defect annealing 2600 nm deep
- Very efficient process Technology applications
- Transferable to other systems

Germanium

Niche applications: Opto-electronics

Purdue – Ultra pure Ge

Low impurities

Ohmic contact – easy (Au-Sb)



Au SBD's

- Resistive evaporation No defects introduced
- good current voltage characteristics

Electrical Characterization of Defects

- Classic DLTS measurements use majority carrier capture and emission process to obtain a "fingerprint" of the defect.
- By monitoring the change of emission rate with temperature an activation energy is obtained.
- By observing the capture rate a cross section can be obtained.



The DLTS Spectrum¹





EMPERATURE

Typical rates: 0.1 - 1000 /s

¹ D. V. Lang, JAP 45, 2023 (1974) or Schroder, "Semiconductor Material and Device Characterization"

Deep Level Transient Spectroscopy

- DLTS requires a depletion (space charge) region
 - p-n, metal-semiconductor, MOS,
- Apply zero / forward bias reverse bias pulse sequence
 - Monitor C-t, I-t, ... as function of T.
 - Transient behaviour indicates the
 Presence of defects with levels
 in the band gap
- Analysis of transient yields:
 - Distiction between majority and minority carrier defects
 - Activation enthalpy (level position), *E*_T
 - Capture cross section, σ
 - Defect concentration, N_T(x)
 - Electric field: defect type (donor, acceptor, ...)
 - Uniaxial stress: defect orientation and symmetry



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L-DLTS of $(H_{0.26} + H_{0.30})$ peak



¹ V. P. Markevich et al, J. Appl. Phys. 95, 4078 (2004).



DLTS: different metallization methods

- Resistive deposition does not introduce defects
 - Curve (a)
- Sputter deposition introduced five electron traps¹
 - Curve (b)
- EBD deposition introduced five electron traps
 - Curve (c)



All the SD induced defects are also introduced by MeV electron irradiation

Curve (d)

¹ F. D. Auret *et al*, J. Electron. Mat. **36**, 1604 (2007).

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DLTS: Example, two RW: Defect E center



$$RW_1 = 80s^{-1} = e_n(T_1) =$$
$$\sigma_n \gamma_n T_1^2 \exp(-\frac{E_T}{kT_1})$$

$$RW_2 = 200s^{-1} = e_n(T_2) =$$
$$\sigma_n \gamma_n T_2^2 \exp(-\frac{E_T}{kT_2})$$

 σ_n and E_T can be determined

Number of traps for n-type semiconductor:



 N_D :concentration of donors

$$\frac{\Delta C}{C_0} = \frac{N_t}{2N_D}$$

The number of traps N_t can be determined

The depletion layer with W increases with the bias potential:

The profile $N_t(\mathbf{x})$ can be obtained

Annealing at room temperature



$$\frac{dN_t}{dt} = -k(T)N_t$$

$$k(T) \ \alpha \exp(-\frac{E_0}{k_B T})$$

Annealing rate constant k(T) can be determined

Isochronal annealing: same time, different temperatures

Annealing activation energy E_0 can be determined

Sample Preparation

Chemical cleaning, ie degreasing
 Chemical etch – oxide removal
 RF Sputter etch – Ar ICP





Use COPRA plasma source

- Inductively coupled plasma (ICP)
- Low energy Ar ions: 1 120 eV
- Fluence rate: $\pm 10^{15}$ cm⁻² s⁻¹
- Etch rates: \pm 0.1 nm s⁻¹ for Ge
- Area: several tens of square cm



4 eV Ar ICP: Average at sample 3.7 eV transfer to Ge atom Sample temperature increases 3 x 10 minute ICP (40°C) Annealing also in time and at increased temperature Sample temperature increases 3 x 10 minute ICP (40°C) Hydrogen passivation of defects

Our basic experiment: 4 eV Ar-ICP plasma



1.-Sb doped Ge is damaged with 5 MeV alpha particles Rest - 24 hours 2.-Au diode is evaporated in half the sample (half A) 3.- DLTS on A (black) 4.-ICP on A and B: 3 x 10 min 5.- Au RE - diode B 6.- DLTS on A (red-dashed) 7.- DLTS on B (blue)



4 eV ICP: E-center concentration -30% Effect depth in Ge: exceeds 2600 nm ICP through metal – less annealing Annealing diminished with higher sample temperature. 1 x 30 minute ICP (70°C) 8 eV Ar ICP (T?) Compare to anneal with phonons: ~ 150 °C Low T (40 °C)

High efficiency

~ 150 °C Lower efficiency Our hypothesis: Ar ions impacting on Ge produce Intrinsic Localised Modes that travel through Ge and anneal defects. Why?

- 1.-ILMs with MD in metals by Hyzhnyakov group have 0.5-5 eV
- 2.-The maximum energy transfer from Ar to Ge is 3.6 eV
- 3.- The activation energy for annealing an E center is about 1.36 eV
- 4.- Energy remains localized exceeding 4000 lattice units
- 5.- Increasing the energy of the plasma does not enhance the effect, this is because ILMs typically have a definite range or energies.
- 6. Increasing sample temperature effect diminished.
- 7.- At least stationary ILMs have been obtained for Si and Ge with MD.

Conclusions:

Plasma of 4eV produces annealing of defects very deep in Ge.

The energy delivered to E-center is ~1 eV.

Likely conclusions:

1. 4 eV Ar hit produces an ILM in Ge with very high efficiency .

 ILM of energy ~3eV travel distances of the order of 10⁴ lattice units or more.

3. The annealing efficiency of ILM with respect to phonons is extremely large.



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Our basic experiment: Facts

1.-Sb concentration: $1.3 \cdot 10^{15}$ cm⁻³ (n_i= 2.4 \cdot 10^{13} cm⁻³) ; 1 Sb per 10⁸ Ge

- 2.- Metal (Au) thickness: 25nm
- 3.- After ICP on Ge the E center concentration drops 29% from N_T = 1.07.10¹⁴ cm⁻³
- 4.- If ICP is done on Au, the E center reduction is smaller, but exists.
- 5.- ICP is done for 30' in intervals to prevent heating
- 6.- Defect annealing occurs up to 2600 nm or 4600 lattice units
- 7.- If the plasma energy is increased the effect is smaller
- 8.- Thermal annealing has to be done at 150°C to obtain a similar effect.

Some numbers

 $\Phi_{Ar} = 1.25 \cdot 10^{13} \text{ cm}^{-2} \text{s}^{-1} \approx 0.04 a^{-2} s^{-1}$ 1.-Ion current can be measured, $\gamma: \quad \Phi_{DB} = \gamma \cdot \Phi_{Ar} \quad ; \quad \gamma < 1$ 2.- DB creation efficiency: $n_{DB} = \frac{\Phi_{DB}}{c_s} \approx \gamma 2.3 \cdot 10^7 \,\mathrm{cm}^{-3}$ 3.-Number of breathers: 4. – Phonons: $E_{ph} = 0.035 \text{eV}$ $n_{ph} = \frac{3n_{Ge}}{\exp(E_{ph}/k_BT) - 1} \approx 4.6 \cdot 10^{22} \text{ cm}^{-3}$ For $\gamma = 1$ $-\frac{1}{N_T} \frac{dN_{T,DB}}{dt} \approx 2 \cdot 10^{-4} s^{-1}$ $-\frac{1}{N_T} \frac{dN_{T,RT}}{dt} \approx 2 \cdot 10^{-11} s^{-1}$

Relative annealing efficiency per DB or phonon: ~10²²

Relative annealing efficiency per eV of DB or phonons: ~10²⁰

Interaction cross-section and energy delivered by a breather

Interaction cross-section σ

$$-\frac{dN_T}{dt} = \sigma N_T \Phi_{DB} \exp(-\frac{E_0 - \Delta}{k_B T})$$

Interaction cross-section

$$\sigma = \alpha \sigma_0$$
 ; $\alpha > 1$

Minimal interacion cross-section

 $\sigma_0 = n_{Ge}^{-2/3} \approx 10^{-15} \text{ cm}^2$

Apparent diminution of the activation energy because of DB interaction: Δ

$$-\frac{dN_T}{dt} = \alpha \gamma \sigma_0 N_T \Phi_{ions} \exp(-\frac{E_0 - \Delta}{k_B T}) \qquad \alpha, \gamma, \Delta \quad \text{unknown}$$
$$-\frac{dN_T}{dt} = \sigma_{ions} N_T \Phi_{ions} \qquad \sigma_{ions} \approx \frac{1}{50} \sigma_0 \approx 1.5 \cdot 10^{-17} \text{ cm}^2$$

For $\alpha \gamma = 1$, $\Delta \approx 1.2 \text{eV}$

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Electron beam deposition

Advantages:

- Easily evaporates high melting point metals.
- Highly controllable deposition rates.
- Good adhesion.

Disadvantage:

- Introduces defects at and below the surface of semiconductors.
- Chen et al
- Mooney et al



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E vs AMU for e or H



 11.5 eV to create Frenkel pair
 1st: electron-atom
 2nd atom to crystal
 Example: e to H to Ge

The Defect Concentration

Defect concentration from peak height, ∆C/C:¹



$\square N_T$ as low as 10¹⁰ defects /cm³

¹ Y. Zohta et al, JAP. 53, 1809 (1982)

sergio@up.ac.za

Depth Distribution of EBD defects

- Use fixed bias, variable pulse DLTS method
 - With " λ " correction¹.
- Not possible to profile the hole traps:
 - Hole concentration is not known.
- V-Sb (E_{0.38}):
 - Concentration decreases rapidly away from the surface.
 - Approaches 10¹⁴ cm⁻³ at surface.
 - Diffusion of vacancies from the surface?



E-beam deposition - Pt



10 keV electrons Beam path Reflected electrons Particles

Experiment 1

E-beam shielding e-traps



High vacuum-2 shields F/gas – 2 shields F/gas – 1 shield F/gas – no shields "Standard EBD"

E-beam shielding h-traps



High vacuum-2 shields F/gas – 2 shields F/gas – 1 shield F/gas – no shields "Standard EBD"





Conclusions 1



Experiment 2 - EBE Deposition vs exposure



No similarities

Arrhenius plots – L-DLTS





New defects – impurity related?

Experiment 2 Au & e-traps



Arrhenius plots – L-DLTS



E_{0.38} Common to both

New defects – perhaps not impurity related

Summary and Conclusions

 1. Shielding lowers the defects introduced by EBD. Damage caused by impacts with ions / particles – not e.
 2. Damage caused in 1st 0.5 µm, at / near defect site. How was the energy transferred?
 3. Damage caused by E < 1.3 eV transfer to Ge. Only enough to displace H or light atoms – single bond.

Will DLTS be more useful in future?



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MeV electrons



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Electron beam deposition



E vs AMU for e or H



I-V: Pt diodes



Ideality = 1.02

Previous lowest: 1.05

Current – lowest measured

Defects linked to e-beam

Next experiment

Defects in Semiconductors

Defects can be "good" or "bad"

- Solar cells: "bad": eliminate them!!!
- Fast switches: "good": deliberately introduce them!

Defects are introduced during

- Crystal growth, sawing / cutting and polishing
- Critical processing steps
 - » Surface cleaning by particle processing (sputter etching)
 - » Metallization
- Radiation
 - » Space, reactors
 - » Accelerators / implanters

Important defect parameters

- Energy level, *E*_T, in bandgap
- Capture cross section, σ
- Concentration, N_T

